



Numerical simulation of a droplet icing on a cold surface

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Outline

1. Introduction

- 2. Physical model and experiments
- 3. Numerical method
- 4. 1D problem and validation
- 5. Simulations of the droplet
- 6. Conclusions

Introduction



Effects

- Lift variation \rightarrow stall
- Increased drag
- Instrumentation problems \rightarrow AF447
- Increased weight

Countermeasures

- De-icing & anti-icing fluids
- De-icing boots
- Electrical resistances
- Hot air bleeding from engines



Physical model and experiments

Freezing droplet



Aim of the work

- Freezing front evolution
- Final shape of the droplet and its causes
 - 1. Density variation
 - 2. Marangoni effect



• solid phase:
$$0 \leq x < s(t)$$

$$\frac{\partial T_s}{\partial t} = \alpha_s \frac{\partial^2 T_s}{\partial x_s^2}$$

• liquid phase: $s(t) < x < \infty$

$$\frac{\partial T_l}{\partial t} = \alpha_l \frac{\partial^2 T_l}{\partial x_l^2}$$

Boundary conditions

- Constant temperatures $T\left(0,t
 ight)=T_{w}$ and $T\left(x
 ightarrow\infty,t
 ight)=T_{i}$
- Interface condition at x = s(t):

$$\begin{cases} -\lambda_{I} \frac{\partial T_{I}}{\partial x} \Big|_{s^{+}} + \lambda_{s} \frac{\partial T_{s}}{\partial x} \Big|_{s^{-}} = \rho_{s} L \frac{ds}{dt} \\ T_{s} \Big|_{s^{-}} = T_{I} \Big|_{s^{+}} = T_{m} \end{cases}$$

Assumptions

- Heat transfer is driven by the sole conduction
- Thermophysical properties constant with temperature in each phase
- Phase change temperature fixed and known
- The entire domain initially at $T(x, 0) = T_i$

Analytical solution

$$s=2\delta\sqrt{\alpha_s t}$$

$$T_{s}(x,t) = T_{wall} + \frac{(T_{m} - T_{wall})}{erf(\delta)} erf\left(\frac{x}{2\sqrt{\alpha_{s}t}}\right)$$
$$T_{l}(x,t) = T_{i} - \frac{(T_{i} - T_{m})}{erfc(\delta\alpha)} erfc\left(\frac{x}{2\sqrt{\alpha_{l}t}}\right)$$

Adimensional parameters

$$Ste = \frac{\rho_s c_{p,s} \left(T_m - T_{wall} \right)}{\rho_s L}$$
$$\phi = \frac{\rho_l c_{p,l} \left(T_i - T_m \right)}{\rho_s c_{p,s} \left(T_m - T_{wall} \right)}$$
$$\alpha = \sqrt{\frac{\alpha_s}{\alpha_l}}$$



Equation for δ

$$\frac{e^{\delta^2}}{erf(\delta)} - \frac{e^{-\delta^2\alpha^2}}{erfc(\delta\alpha)}\frac{\phi}{\alpha} = \frac{\delta\sqrt{\pi}}{Ste}$$

Velocity calculation

Continuity

$$\rho_{liq}V_{liq} + \rho_{sol}V_{sol} = const \qquad \rightarrow \qquad v_{liq} = \frac{dH}{dt} = \frac{ds}{dt} \cdot \underbrace{\left(1 - \frac{\rho_{sol}}{\rho_{liq}}\right)}_{r}$$

Governing equation $\frac{\partial T_l}{\partial t} + v_{liq} \frac{\partial T_l}{\partial y} = \alpha_l \frac{\partial^2 T_l}{\partial x_l^2}$

Being the parameter r the indicator of the expansion:

$$T_{I} = T_{0} + (T_{m} - T_{0}) \frac{\operatorname{erfc}\left[\alpha\delta\left(\frac{x}{2\delta\sqrt{\alpha_{s}t}} - r\right)\right]}{\operatorname{erfc}\left[\alpha\delta\left(1 - r\right)\right]}$$
$$\frac{e^{-\delta^{2}}}{\operatorname{erf}\left(\delta\right)} - \frac{\phi}{\alpha} \frac{e^{-(\alpha\delta)^{2}}}{\operatorname{erfc}\left[\alpha\delta\left(1 - r\right)\right]} \frac{e^{2r(\alpha\delta)^{2}}}{e^{(r\alpha\delta)^{2}}} = \frac{\delta\sqrt{\pi}}{Ste}$$

Numerical method

Jadim is a research code developed by J. Magnaudet and D. Legendre's team in the Interface group at *Institut de Mécaniques des Fluides de Toulouse*. The code permits to describe in an accurate way physical mechanisms present in multiphasic flows.

- Volume of Fluid formulation is employed
- Thermal and Immersed Boundary Method routine are supported

In the present work the objective is to couple the three of them and, in particular, to develop a thermal based IBM formulation

Solid function fraction

A solid function fraction has been defined as follows, representing the amount of ice for each cell: -

$$\alpha_{ibm,lin} = \tau \cdot \frac{T_{max} - T}{T_{max} - T_{min}}$$
$$\alpha_{ibm,cos} = 0.5 \cdot \tau \cdot \left[\cos \frac{\pi \cdot (T - T_{min})}{T_{max} - T_{min}} + 1 \right]$$



$$\phi = (1 - \tau) \phi_{air} + \tau \left[(1 - \alpha_{ibm}) \phi_{liq} + \alpha_{ibm} \phi_{sol}
ight]$$

Velocity calculation

$$oldsymbol{v}_{\mathit{liq}} = \left(1 - rac{
ho_{\mathit{ice}}}{
ho_{\mathit{water}}}
ight)oldsymbol{V}_{\mathit{front}}$$

Scalar transport equation:

(

$$\frac{\partial \alpha_{ibm}}{\partial t} + \mathbf{V}_{front} \cdot \nabla \alpha_{ibm} = 0$$
$$\mathbf{V}_{front} \cdot \mathbf{n} = \frac{-\frac{\partial \alpha_{ibm}}{\partial T} \frac{\partial T}{\partial t}}{||\nabla \alpha_{ibm}||}$$
$$n_i = \frac{\frac{\partial \alpha_{ibm}}{\partial x_i}}{||\nabla \alpha_{ibm}||}$$

Velocity imposition

$$f = \alpha_{ibm} \frac{U - U^*}{\Delta t}$$

$$\begin{cases} U = 0 & \alpha_{ibm} \ge 0.95 \\ U = v_{liq} & 0 < \alpha_{ibm} < 0.95 \end{cases}$$

Being U^* a predictor velocity without considering the immersed object **Pressure correction** Jadim's SIMPLE algorithm is modified in order to take account of the calculated velocities

The latent heat computation

$$egin{aligned} &(H=c_{p,I}T & ext{Liquid}\ &(H=c_{p,s}T+L & ext{Solid}\ & &rac{\partial
ho H}{\partial t} =
abla \cdot (k
abla T) \end{aligned}$$

The apparent heat capacity method

The source term method

$$\rho c_{p} \frac{\partial T}{\partial t} = \frac{\partial}{\partial x} \left(k \frac{\partial T}{\partial x} \right) - S$$
$$S = \frac{\partial \alpha_{IBM}}{\partial T} \frac{\partial T}{\partial t} \frac{[L + T (c_{p,s} - c_{p,l})]}{c_{p}}$$

$$c_{app} = \frac{dH}{dT}$$

$$c_{app} = c_p + L \frac{d\alpha_{IBM}}{dT}$$

$$\rho c_{app} \frac{\partial T}{\partial t} = \nabla \cdot (k \nabla T)$$

1D problem and validation

	err _{T,max}	err _{T,avg}	err _{int,max}	err _{int,avg}
Linear [-2, 2]	6.1%	5.7%	27%	17%
Linear $[-1,1]$	4.3%	4.1%	14.2%	9.2%
Linear [0, 1]	5.2%	4.2%	4.7%	2.2%
Linear [0, 0.1]	2.8%	2%	3.5%	1%
Cosine [-1,1]	5.2%	4.3%	12.2%	7.6%
Cosine [0, 1]	5.3%	4.1%	4.8%	2.3%

- For narrower solidification ranges, the results are more accurate
- The solidification front is well located if the inferior limit coincides to $0^\circ C$

IBM functions, figures





Freezing front, $\cos[-1,1]$

Freezing front, lin [0, 0.1]



Velocity

	err _{vel,max}	err _{vel,avg}
Lin [-1, 1]	16.5%	8.1%
Lin [0, 1]	23.6%	8.7%
Lin [0, 0.1]	100%	53.8%
Cos [-1, 1]	13.7%	7%
Cos [0, 1]	23.6%	8.9%

- Velocity is generally underestimated
- The chosen range has to be wide enough to properly calculate the velocity



Grid convergence



Figure 1: Temperature, $\cos[-1, 1]$



Figure 2: Velocity, $\cos [-1, 1]$



Figure 3: Interface position, $\cos [-1, 1]$



Latent heat



Figure 5: Apparent capacity method



- Average errors are comparable: $\sim 20\%$
- Source term method tends to become more accurate through time
- Apparent heat capacity worsens with time due to underestimation of the latent heat
- Source term method doesn't guarantee stability without additional loops

Figure 6: Source term method

Simulations of the droplet

Parameters of the simulations



$$\begin{array}{c} \alpha \left[\frac{m^2}{s} \right] & \rho \left[\frac{kg}{m^3} \right] \\ \text{Air} & 2.166 \cdot 10^{-5} & 1.2 \\ \text{Water} & 1.433 \cdot 10^{-7} & 1000 \\ \text{Ice} & 1.176 \cdot 10^{-6} & 917 \end{array}$$

$$\begin{array}{lll} \mathsf{V} & 1.35 \cdot 10^{-13} \ [m^3] \\ \mathsf{R}_{90} & 4 \cdot 10^{-5} \ [m] \\ \mathsf{dx} & 10^{-6} \ [m] \\ \mathsf{dt} & 10^{-7} \ [s] \end{array}$$

- West wall represents the cold plate and it is isothermal
- South is the symmetry axis
- North and east walls are adiabatic

 $T_i = 18^\circ C$

 $T_w -5^\circ C$

$$Bo = rac{\mathbf{g}\cdot \Delta
ho \cdot \mathbf{d}^2}{\gamma} \ll 1$$

The solidification process (1)













- Higher thermal diffusivity of the air causes the droplet outer layer to solidify
- Thermally treated as a mixture between ice and water
- Dynamically considered as liquid

Velocity field



- Inside the solid velocity is zero
- Composition of V_x and V_y returns a vector perpendicular to the interface
- Velocity field in the water comes from the combination of incompressibility and the axis of symmetry

Evolution of the freezing front (1)





- Quasi-linear trend till the outer layer of ice is thin
- Acceleration due to the increased thermal diffusivity

Evolution of the freezing front (2)



- Experimentally the trend seems to be linear differently from the Stefan problem
- An expanded domain guarantees a more accurate result, minimum size depends on the contact angle
- Numerical results confirm the global linear behaviour

Contact angle influence (1)

- Wall temperature
- Initial temperature
- Droplet's volume
- Mesh size







Contact angle influence (2)



- Hydrophobic surfaces tends to delay the ice accretion, commonly employed as constructive materials or coatings
- + We observed no formation of the protrusion in correspondence of a threshold of about $\Theta\simeq 30^\circ$

Contact angle influence (3)



- Results similar to other numerical works
- Quasi-linear behaviour independent from the contact angle, the difference is the total solidification time

It is not a physical case \Rightarrow the droplet's volume and consequentially Bond number are too small to cause an elliptical shape

A qualitative study to investigate if new dynamics appear



- The freezing front shape resemble the spheric cap case and the experimental results
- The pointy protrusion continues to appear





Conclusions

Conclusions

- Successfully coupled VoF, thermics and IBM
- The characteristic pointy tip has been reproduced
- Evolution of the freezing front confirms experimental results over analytical ones
- Parametric study of the contact angle influence

Future developments

- Ameliorate the computation of the latent heat of solidification
- Tracking the freezing front for a better calculation of the velocity
- Experimentally confirm the influence of the contact angle

Thank you for your attention