

# INSTABILITY AND TRANSITION OF THE ROTATING DISK BOUNDARY LAYER OVER HOMOGENIZED TEXTURED SURFACES

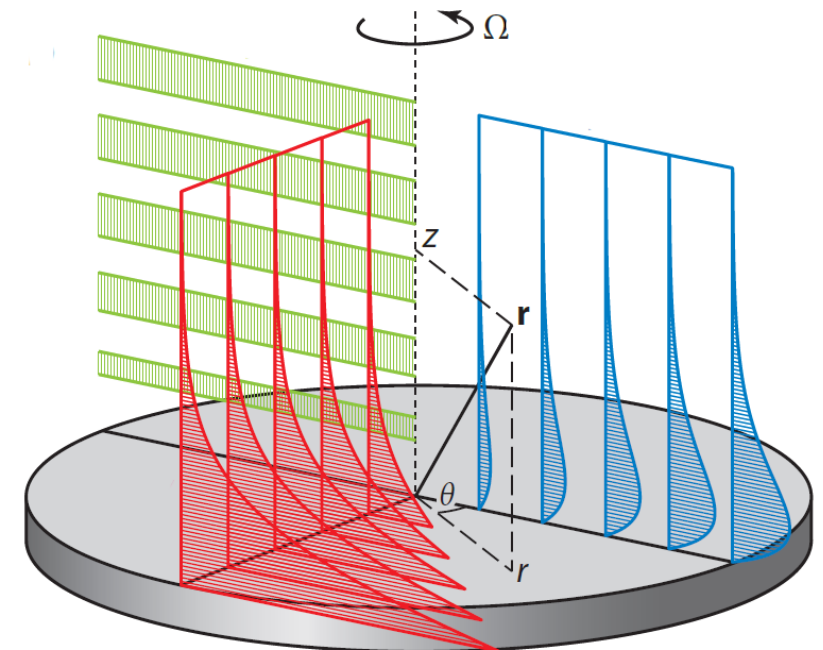
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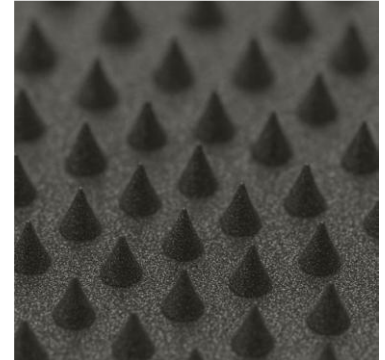


# Textured disks

**Circumferential**

**Radial**

**Isotropic**



# Homogenization theory

	Inner/micro-scales	Outer/macro-scales
Length	$\ell$	$(\nu/\Omega)^{1/2}$
Velocity	$\Omega\ell$	$(\nu\Omega)^{1/2}$
Time	$\Omega^{-1}$	$\Omega^{-1}$
Pressure	$\mu\Omega$	$\mu\Omega$

von Kármán scales

## To set ideas:

Imagine a disk rotating at  $\Omega = 100$  rad/s in water, the macroscopic length scale is 0.1 mm, the laminar boundary layer is about 0.5 mm thick.

The Reynolds number,  $Re = \frac{\hat{r}}{\sqrt{\nu/\Omega}}$ , attains the value of 500 when the radius is  $\hat{r} = 5$  cm.

$\ell$ : characteristic dimension of the wall texture (periodicity of the pattern)

Expand all terms in powers of  $\epsilon = \frac{\ell}{\sqrt{\nu/\Omega}}$  and solve the Stokes equations at different orders of  $\epsilon$  in a unit cell for each of the textured disks considered. Observe that, when  $\epsilon = 1$ , the laminar boundary layer thickness is ten times the height of the roughness.

To first order in  $\epsilon$  the solutions in the unit cells yield the slip lengths,  $\lambda_\theta$  and  $\lambda_R$ , of the **effective wall conditions**

## Effective conditions of the *macroscopic* problem

$$U_R|_{Z=0} = \epsilon \lambda_R S_{RZ}|_{Z=0} + \mathcal{O}(\epsilon^2)$$

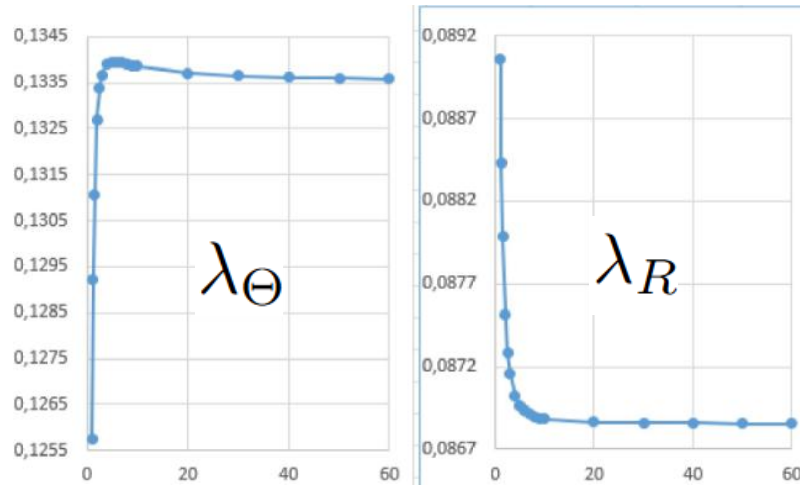
$$U_\Theta|_{Z=0} = R + \epsilon \lambda_\Theta S_{\Theta Z}|_{Z=0} + \mathcal{O}(\epsilon^2)$$

$$U_Z|_{Z=0} = \mathcal{O}(\epsilon^2)$$

$Z = 0$  : **fictitious surface**  
where the effective wall  
conditions are applied

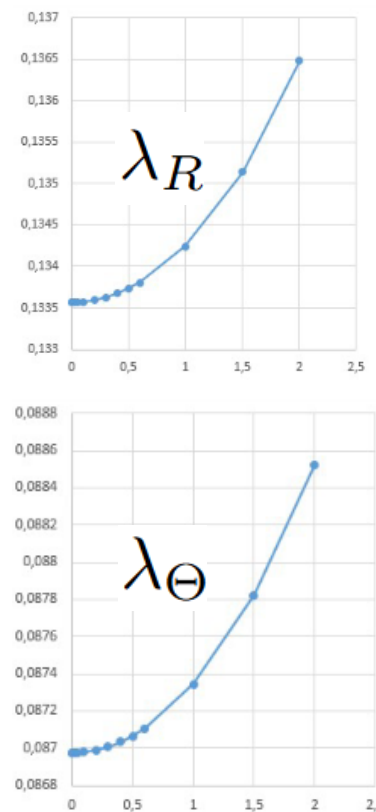
$$S_{RZ} = \frac{\partial U_R}{\partial Z} + \frac{\partial U_Z}{\partial R} \Big|_{Z=0}, \quad S_{\Theta Z} = \frac{\partial U_\Theta}{\partial Z} + \frac{1}{R} \frac{\partial U_Z}{\partial \Theta} \Big|_{Z=0}, \quad S_{ZZ} = -P + 2 \frac{\partial U_Z}{\partial Z} \Big|_{Z=0}$$

## Circumferential ribs



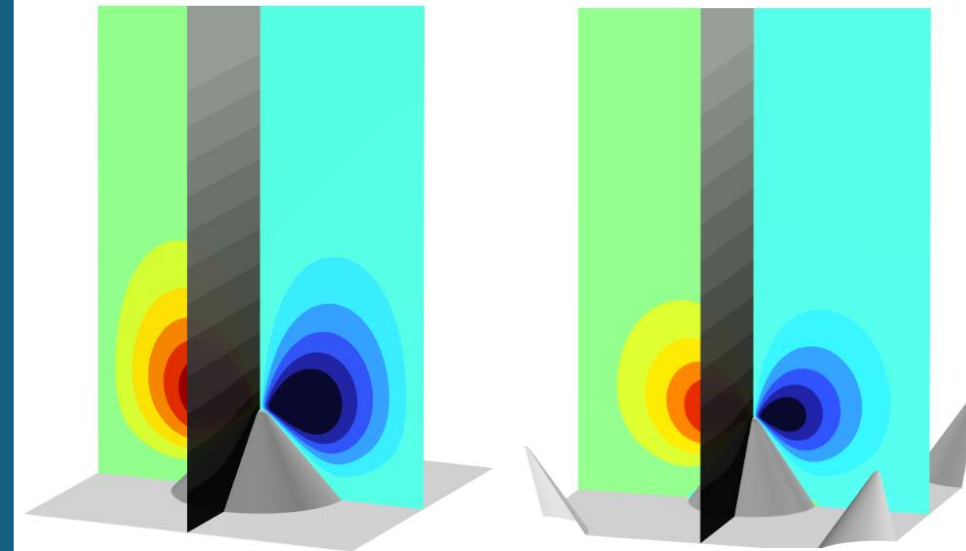
The coefficients vary with the radial position  $\frac{\hat{r}}{\ell}$ . Assuming  $\epsilon = 1$  the radial position at which the slip lengths are independent of the radius is  $R = Re \approx 20$ .

## Radial ribs



They vary with the angular opening of the ribs,  $\Delta\theta$

## Isotropic ribs inline/staggered

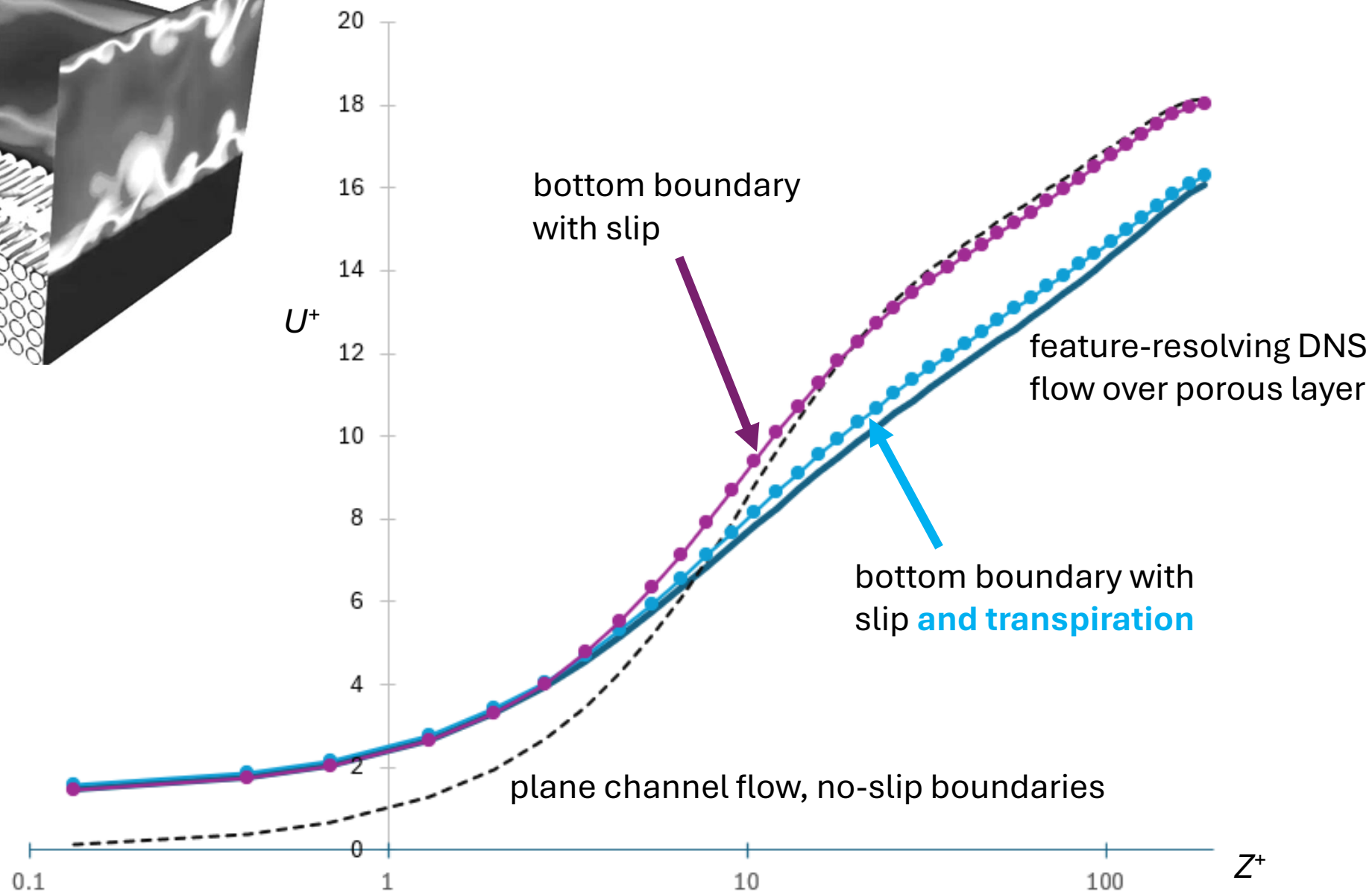
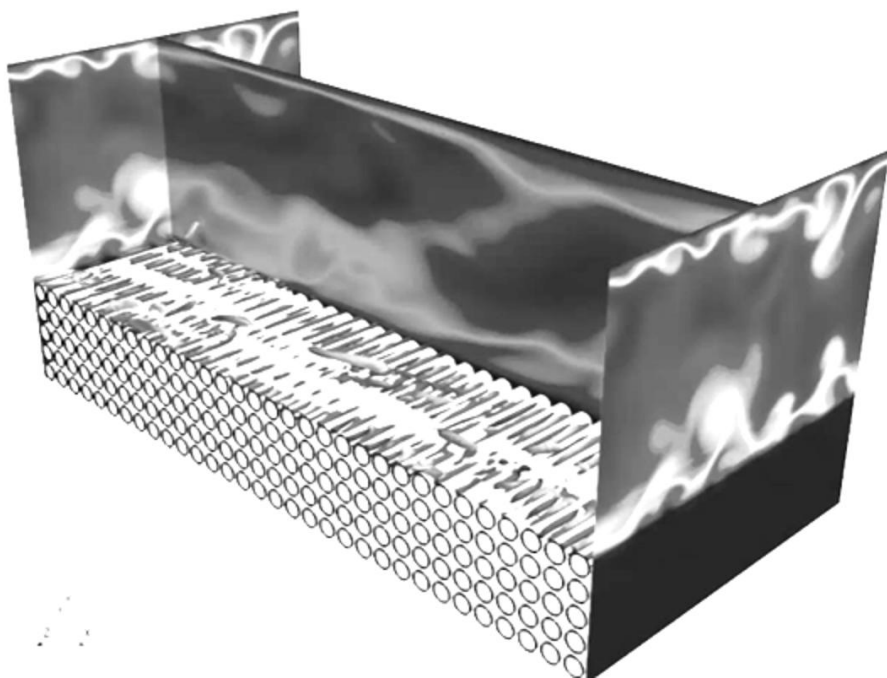


	$\lambda_\Theta$	$\lambda_R$
Conical roughness (inline arrangement)	0.17839	0.17839
Conical roughness (staggered arrangement)	0.15880	0.15880

Are the boundary conditions at leading order in  $\epsilon$  sufficient?

Are the boundary conditions at leading order in  $\epsilon$  sufficient?

The answer is **yes**, if the flow is laminar,  
but **no** after transition to turbulence has occurred





## THE TRANSPIRATION BOUNDARY CONDITION IN THE DISK CASE

$$U_Z|_{Z=0} = -\epsilon^2 \left[ \mathcal{K}_{\Theta Z}^{itf} \frac{\partial}{\partial Z} \left( \frac{1}{R} \frac{\partial U_{\Theta}}{\partial \Theta} + \frac{U_R}{R} \right) + \mathcal{K}_{RZ}^{itf} \frac{\partial^2 U_R}{\partial Z \partial R} \right] \Big|_{Z=0} + \mathcal{O}(\epsilon^3)$$

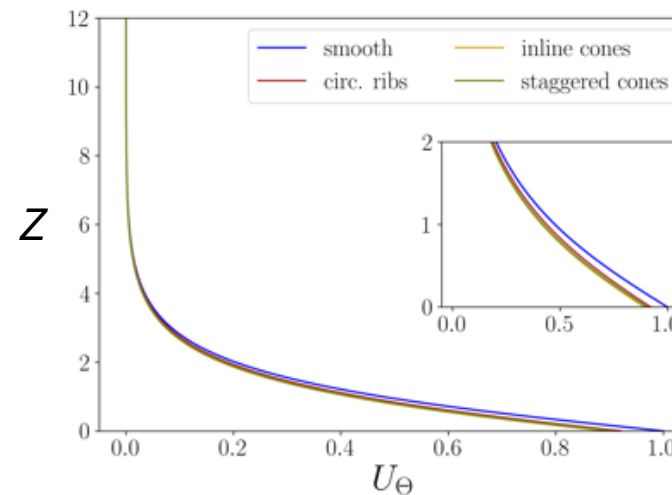
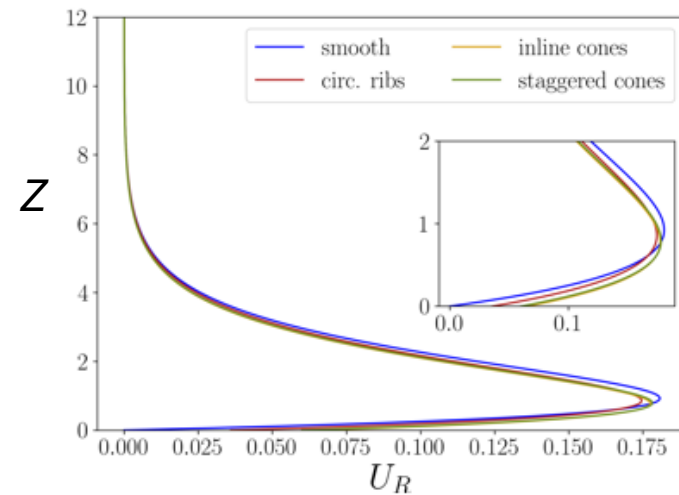
For isotropic roughness elements the condition can be simplified to read

$$U_Z|_{Z=0} = \epsilon \lambda_Z \frac{\partial U_Z}{\partial Z} \Big|_{Z=0} + \mathcal{O}(\epsilon^3),$$

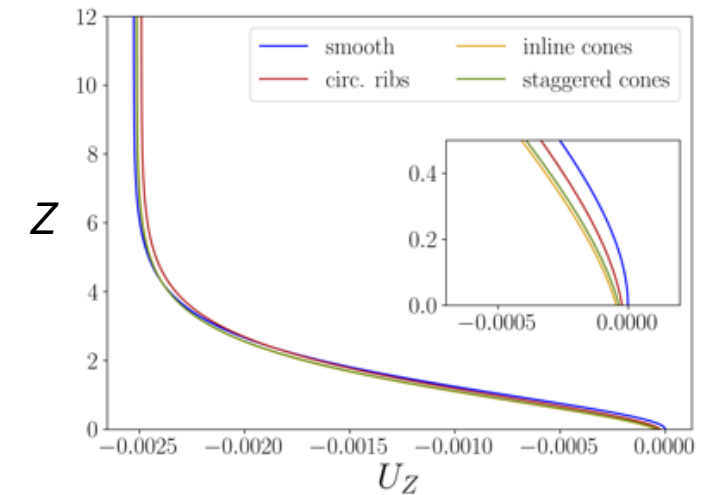
with  $\lambda_Z$  a small, formally  $\mathcal{O}(\epsilon)$ , wall-normal protrusion height.

# Back to the disk ...

**Laminar base flow**  
von Kármán solution  
(with and without  
slip/transpiration)

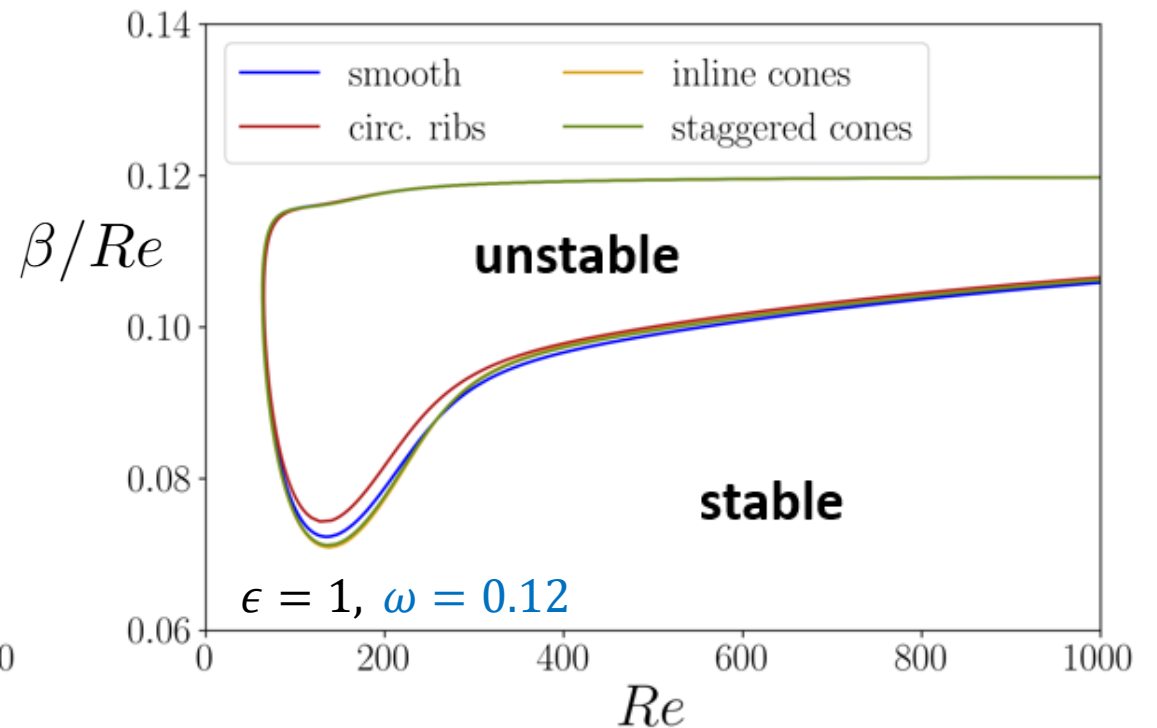
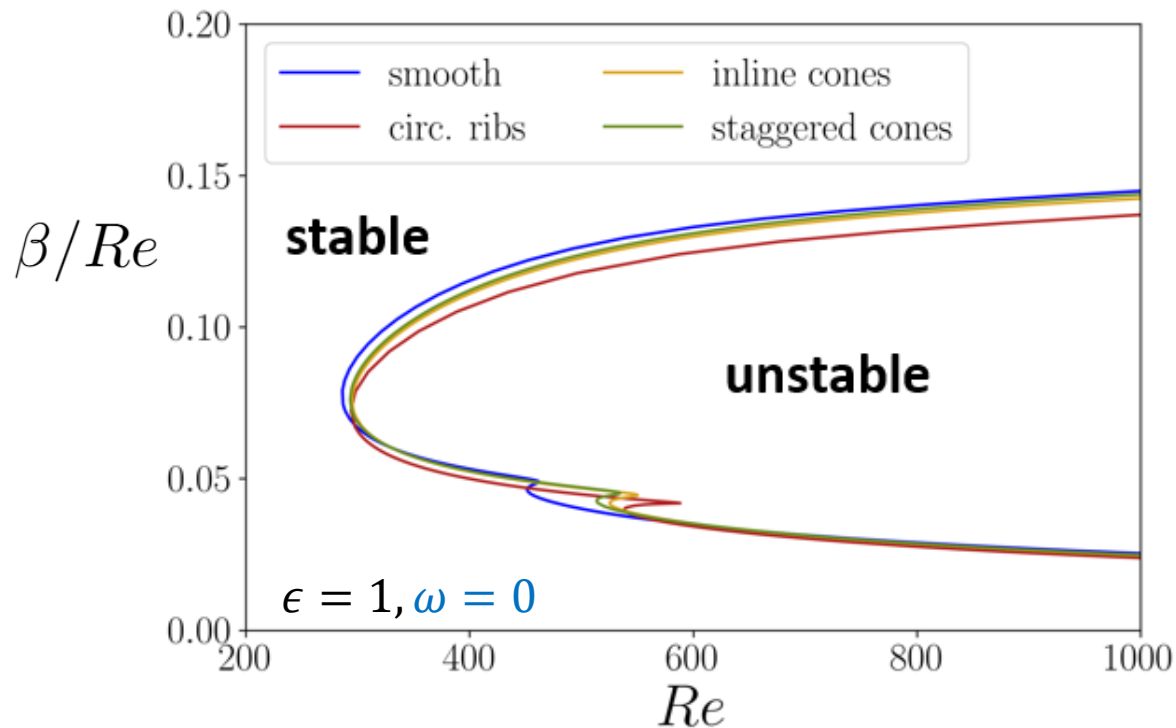


Laminar solution for  $\epsilon = 1$  and  
 $R = 350$



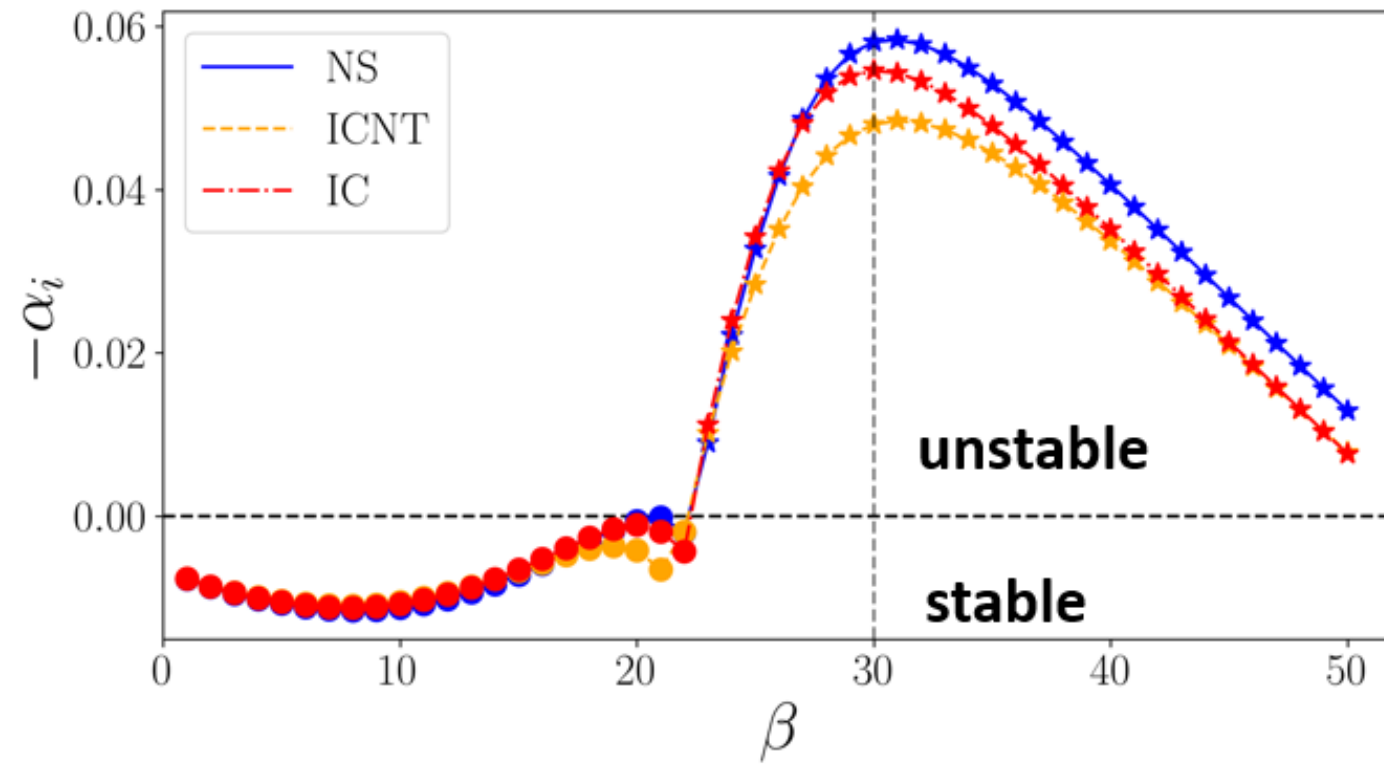
# Local linear stability

Take small disturbances of the form:  $f(\hat{z}) \exp[i(\alpha \hat{r} + \beta \theta + \omega \hat{t})]$  and search for complex eigenvalues  $\alpha$



Neutral curves for textured surfaces with and without transpiration are overlapped

... but at  $Re = 450$  the growth rates differ!

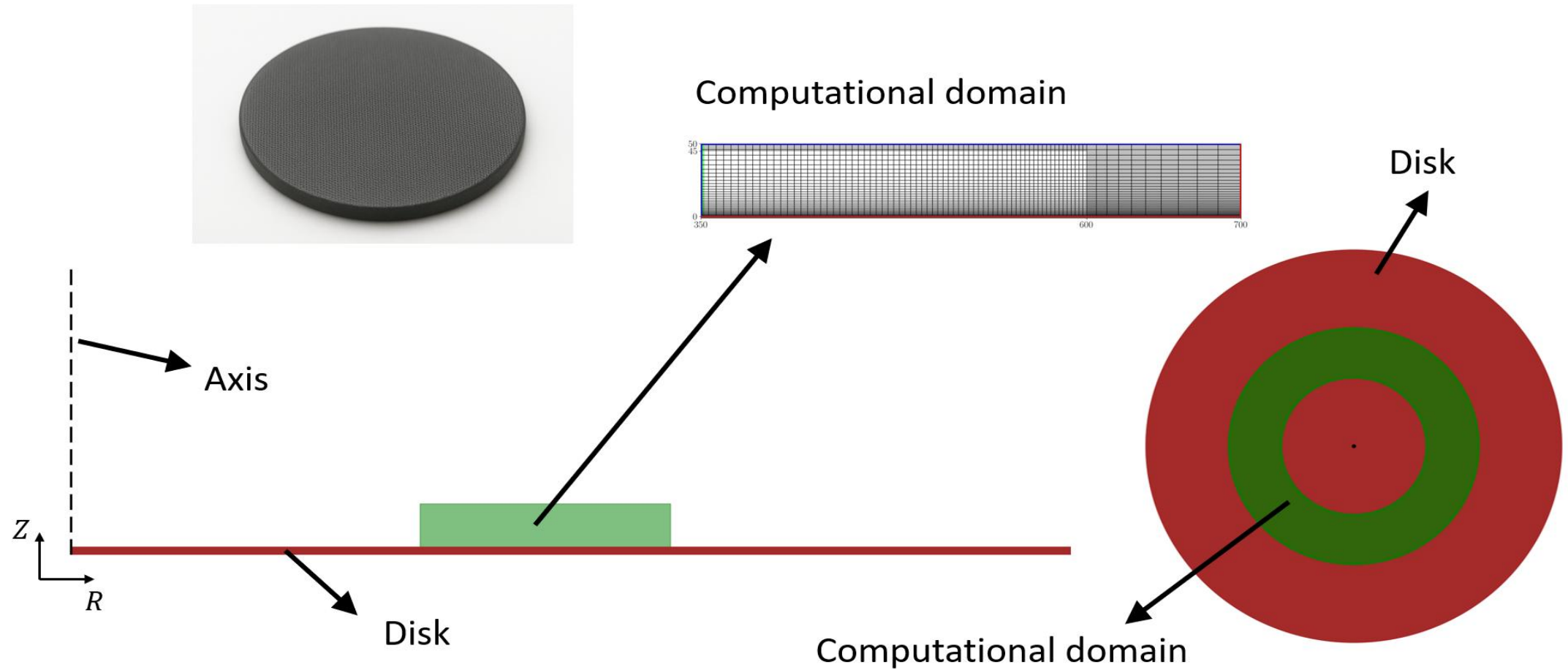


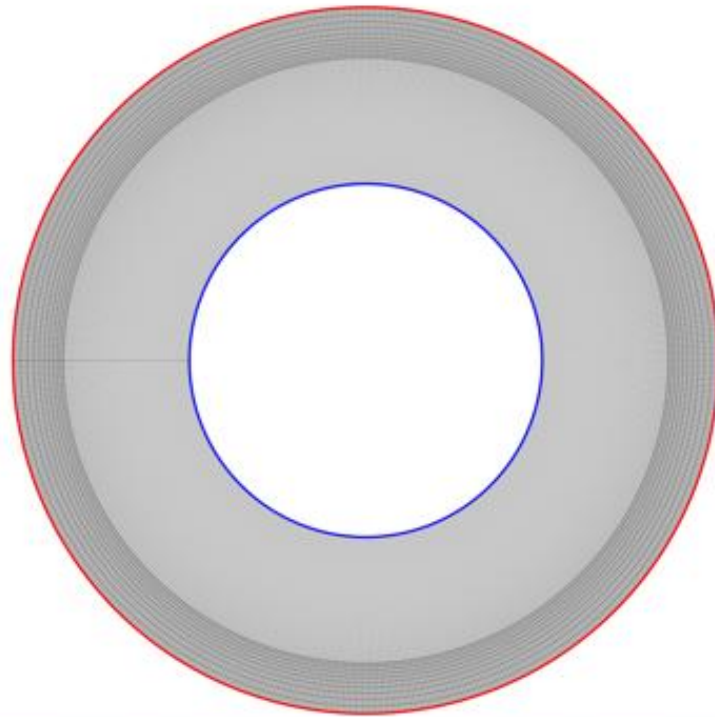
NS = no-slip

ICNT = inline cones, no transpiration

IC = inline cones, w/ transpiration

# DNS

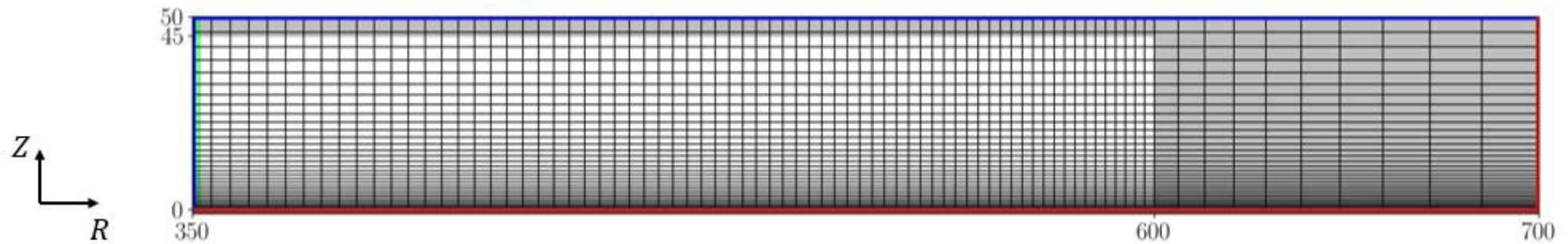


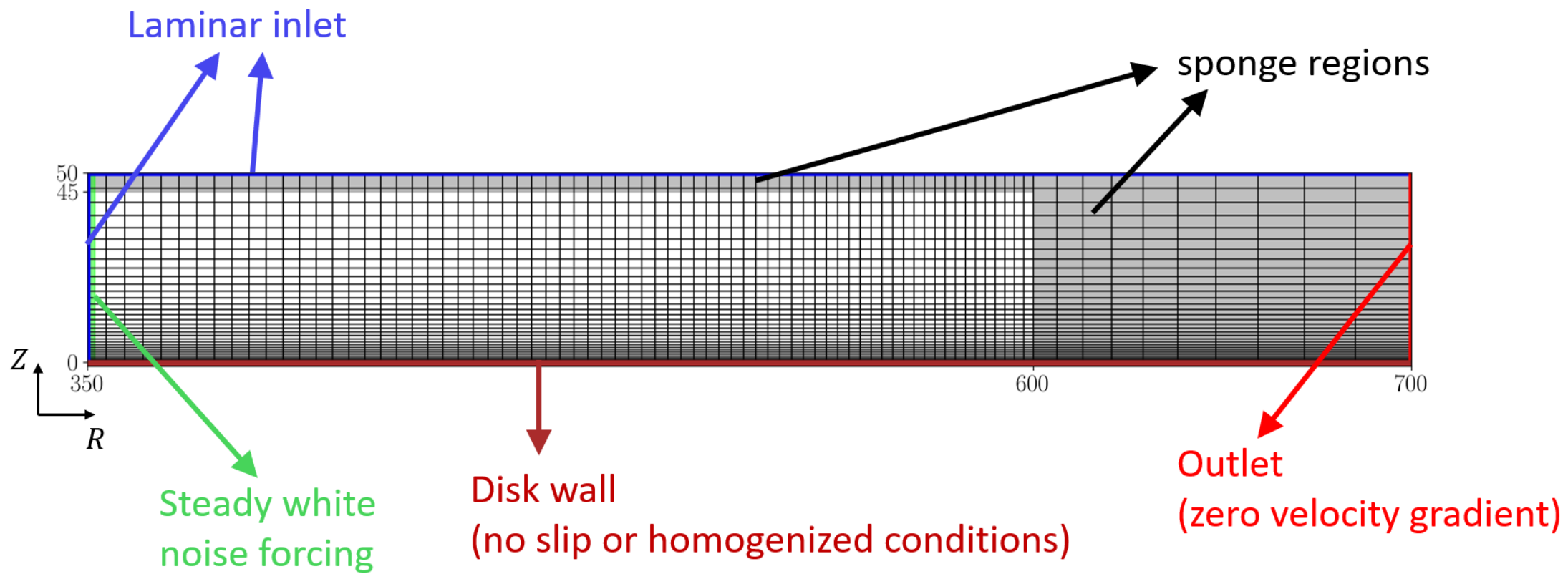


Nek5000 spectral element code  
(7<sup>th</sup> order Legendre polynomials)

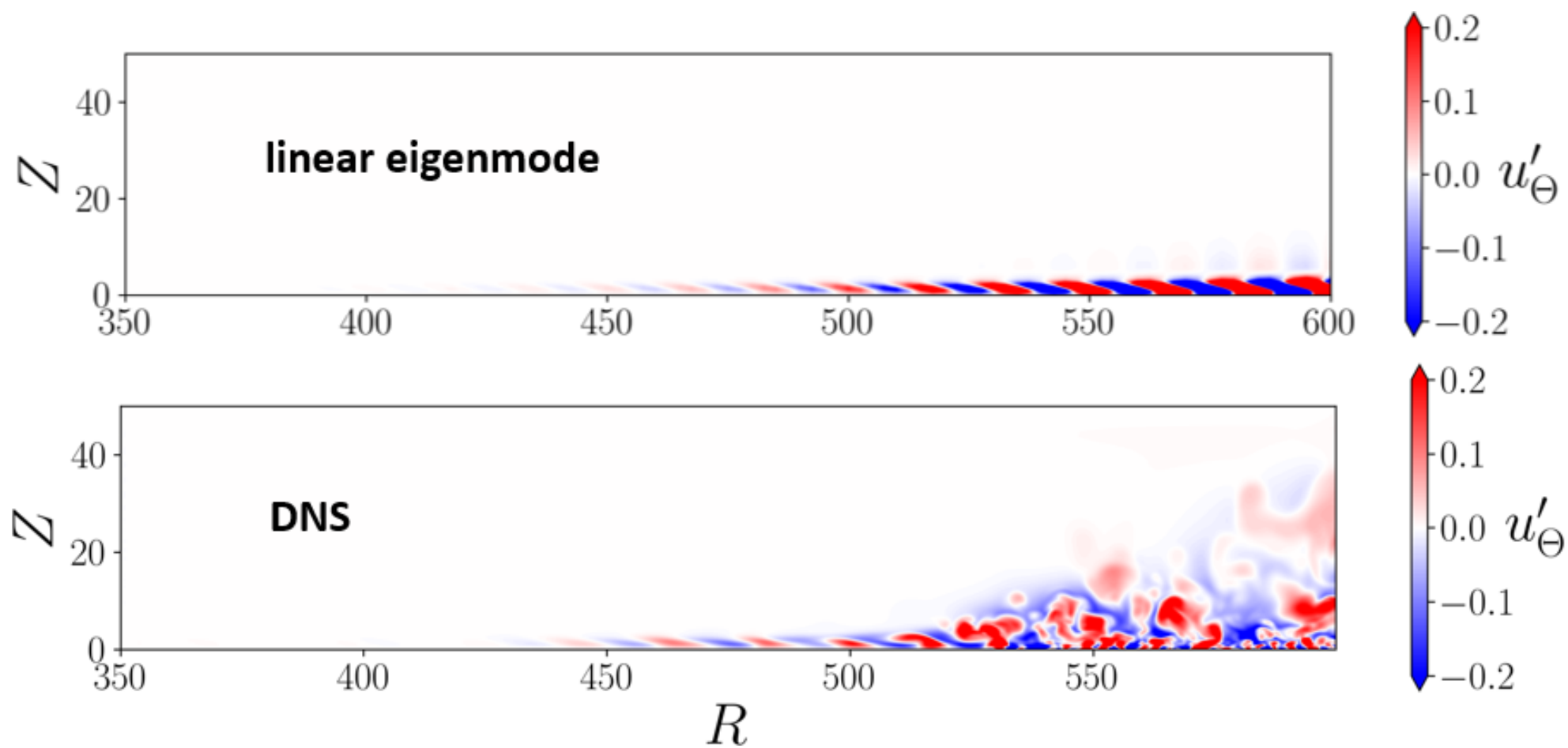
$\approx 2 \cdot 10^6$  spectral elements

$\approx 10^9$  grid points

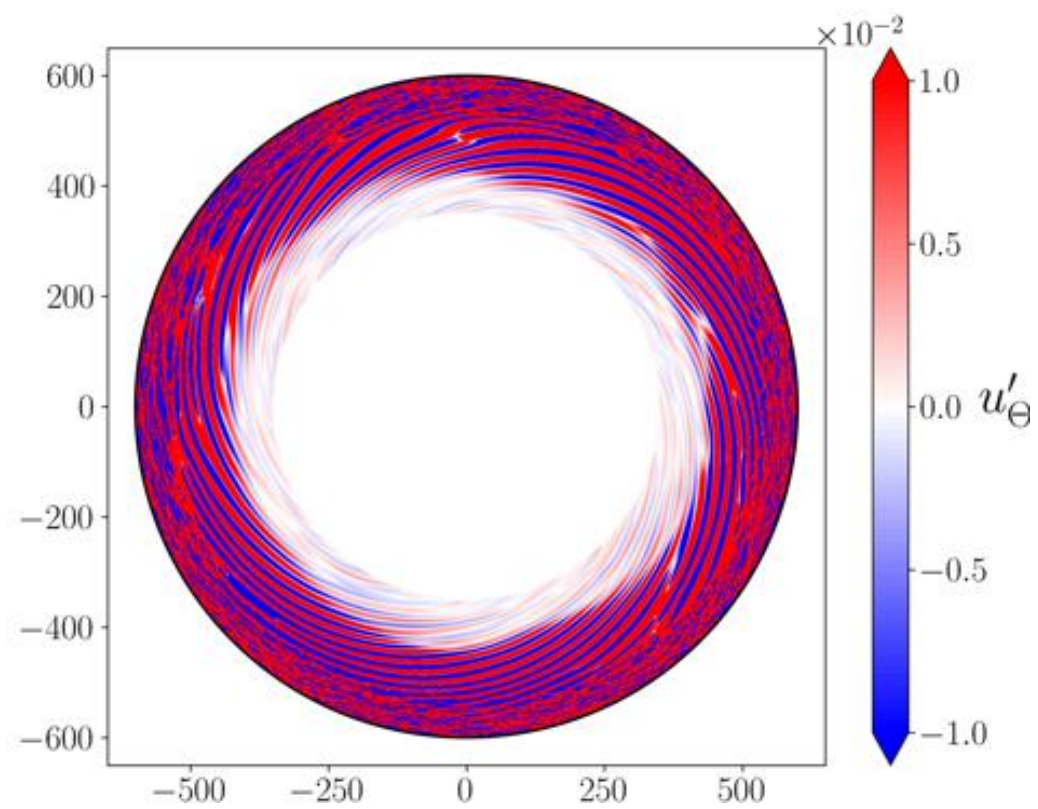
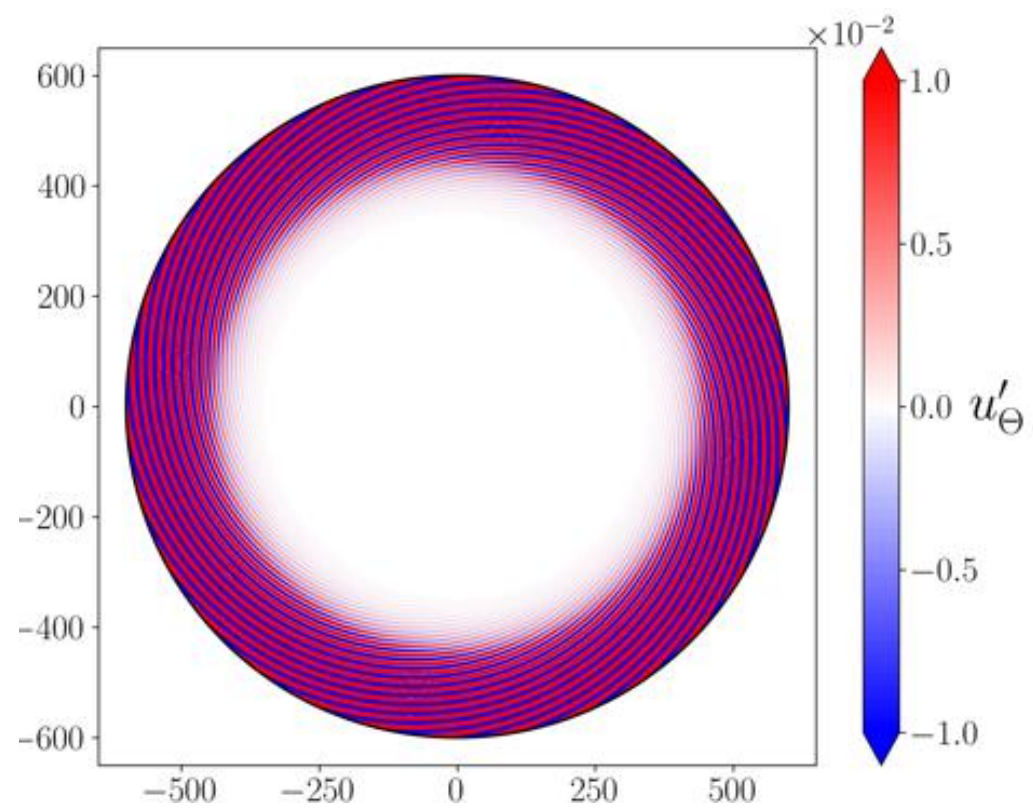




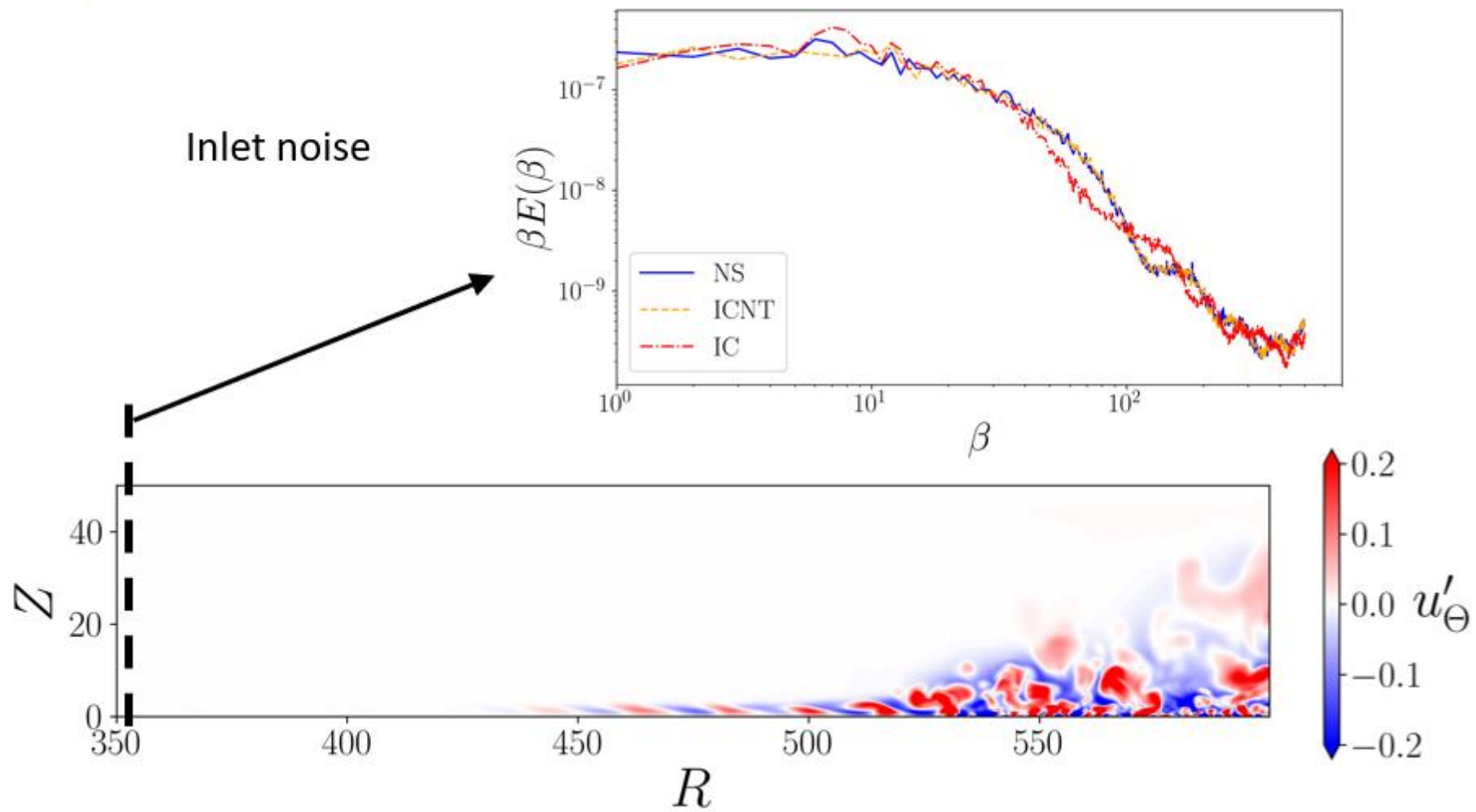
IC case: inline cones, with transpiration

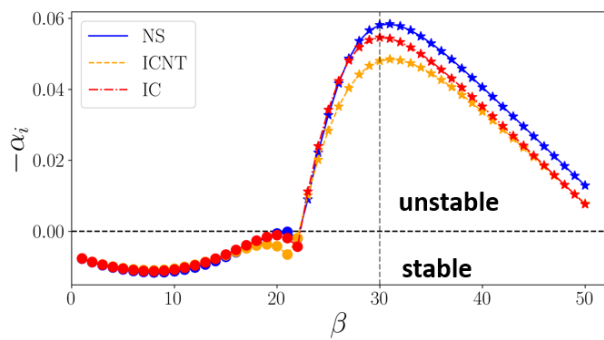




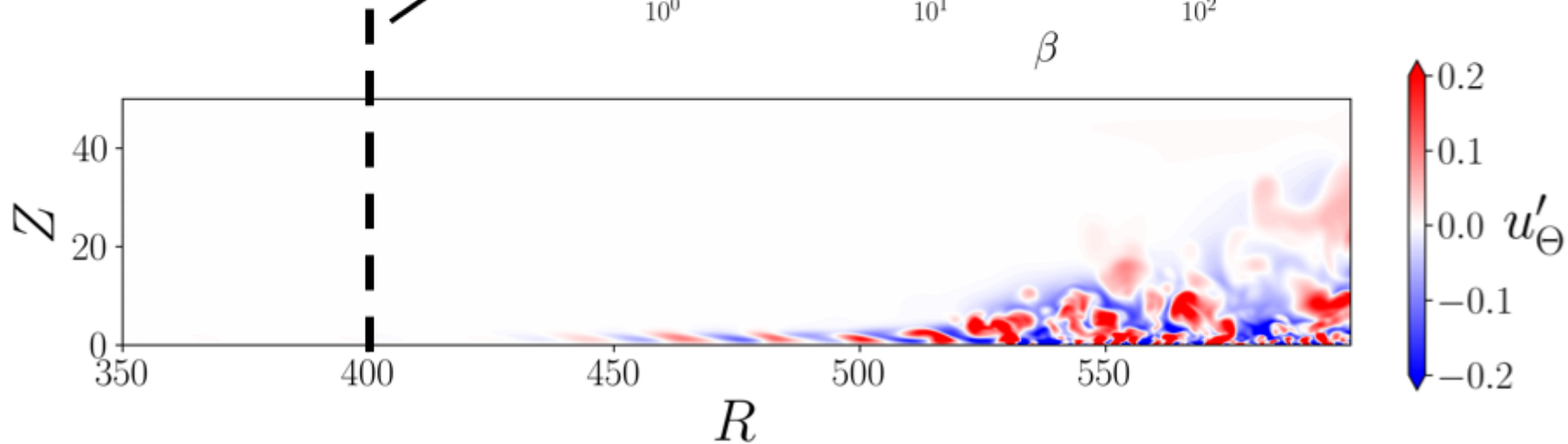
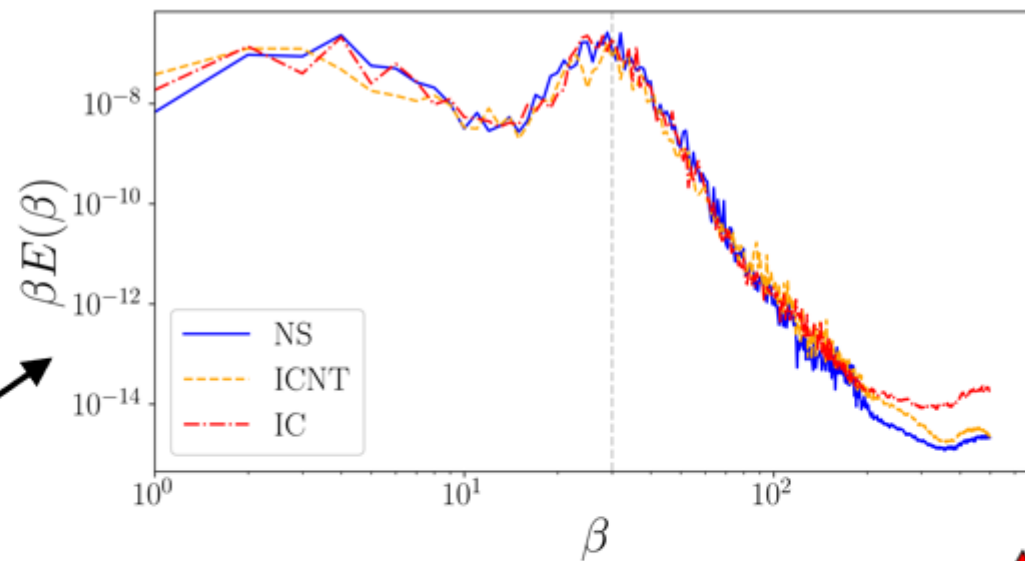


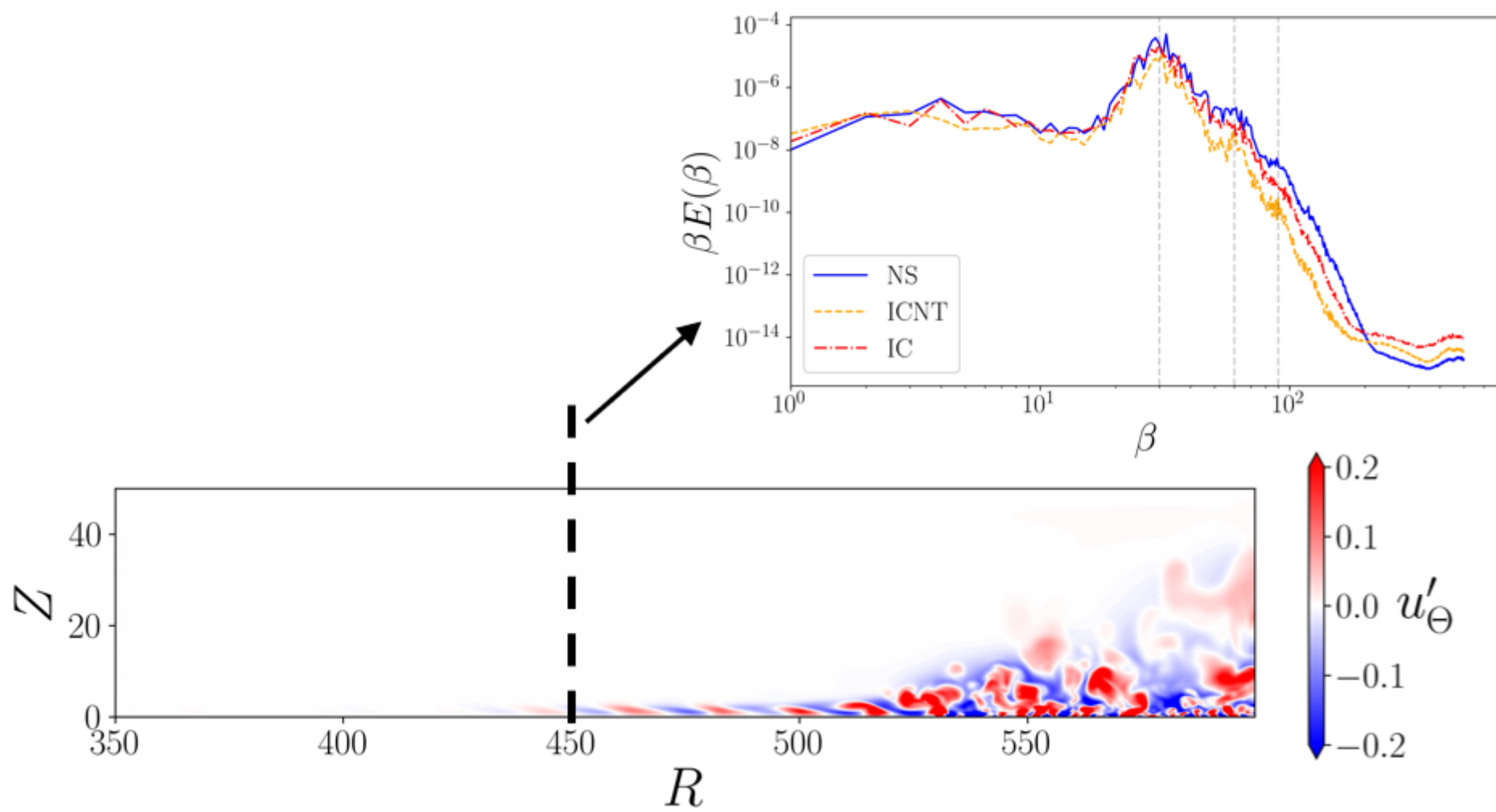
Wall-normal integrated pre-multiplied  
azimuthal Fourier spectrum

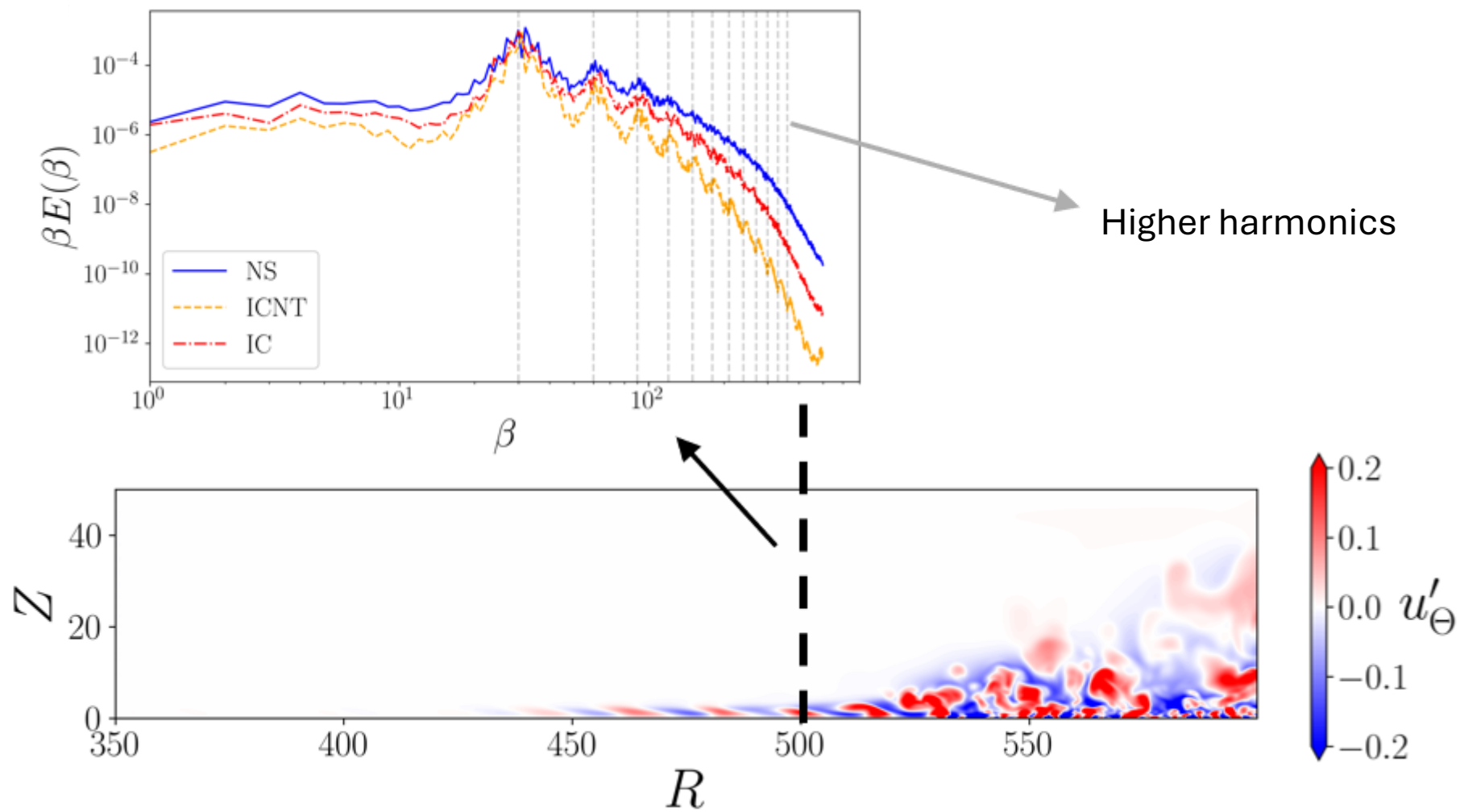


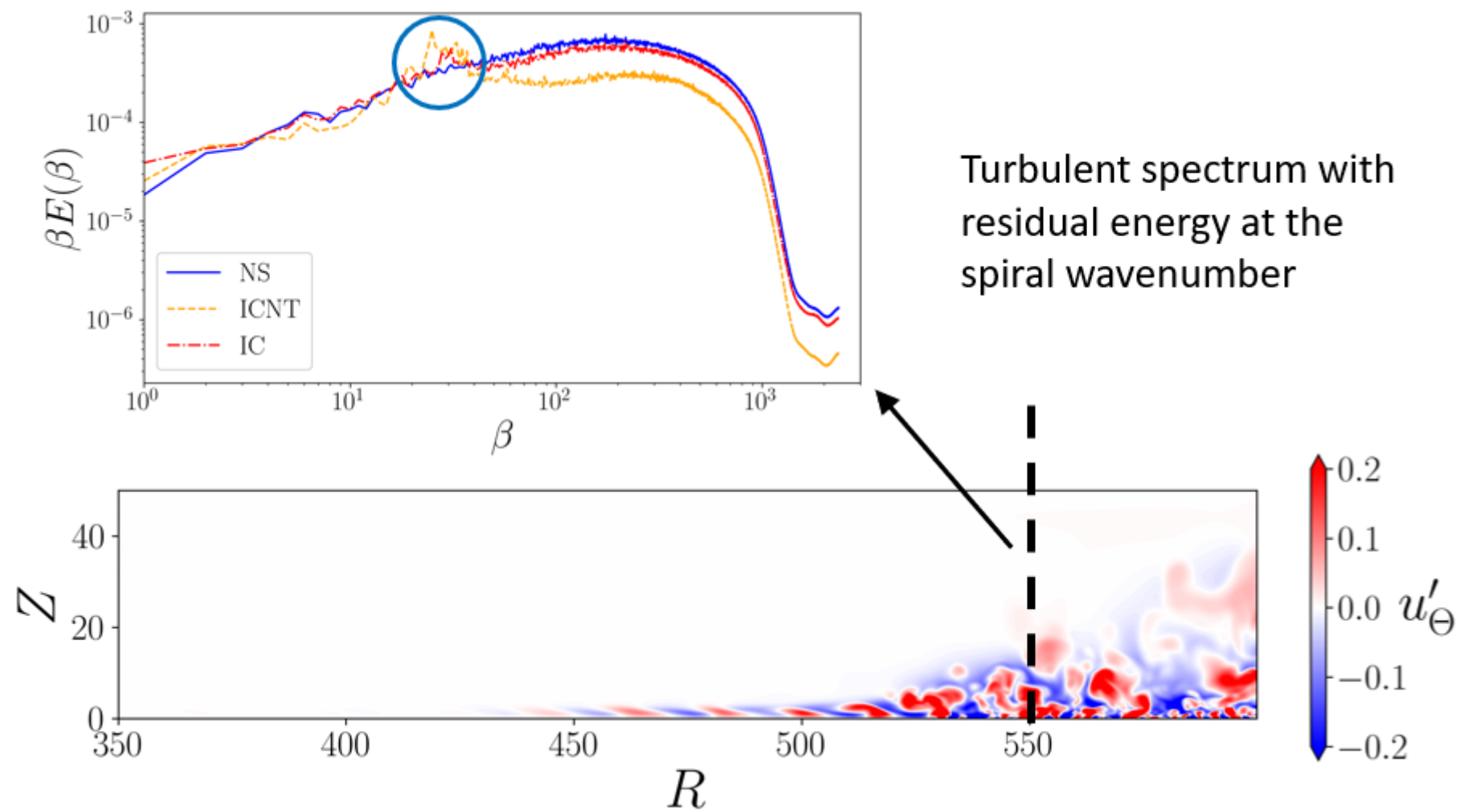


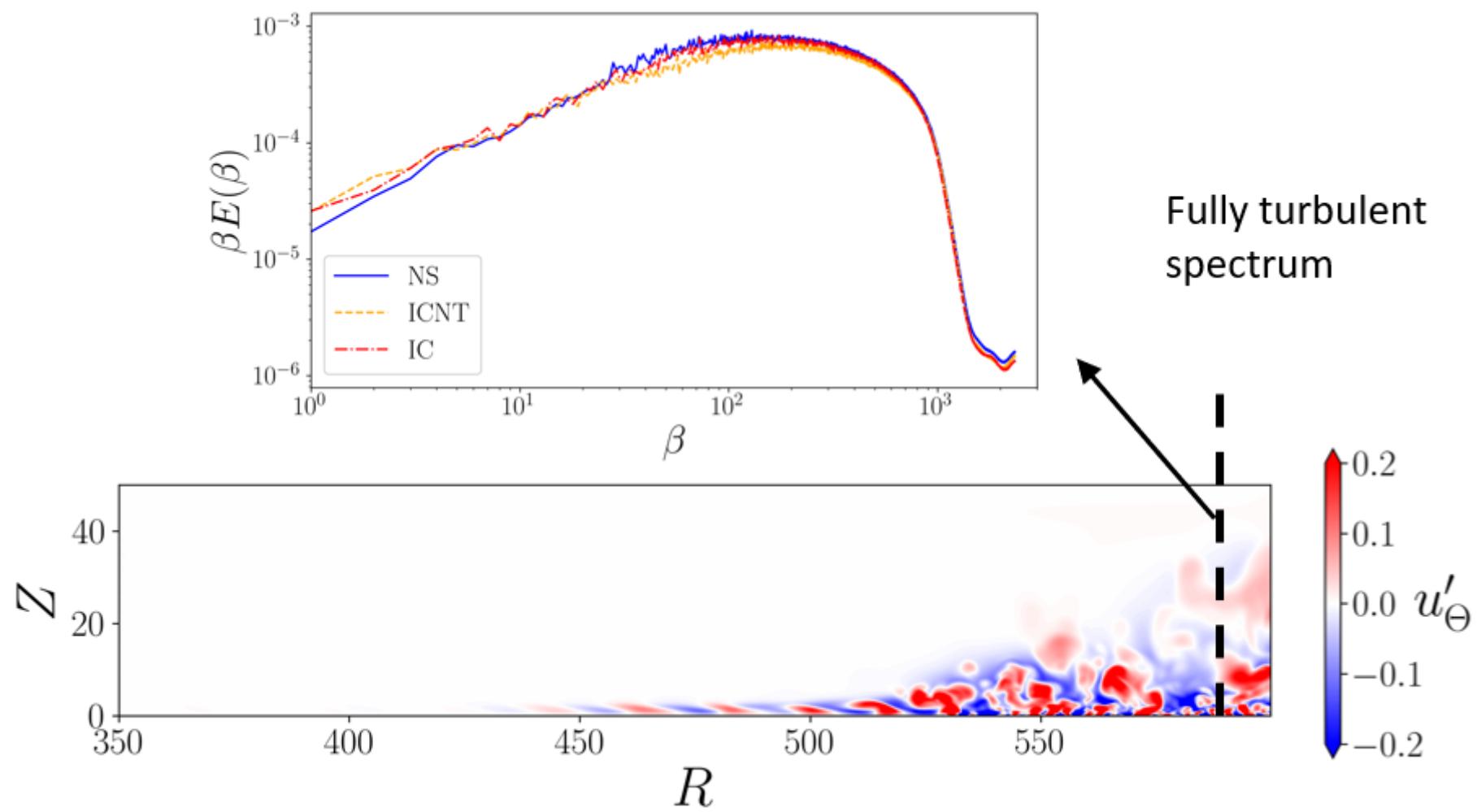
Emergence of  
coherent  
structures

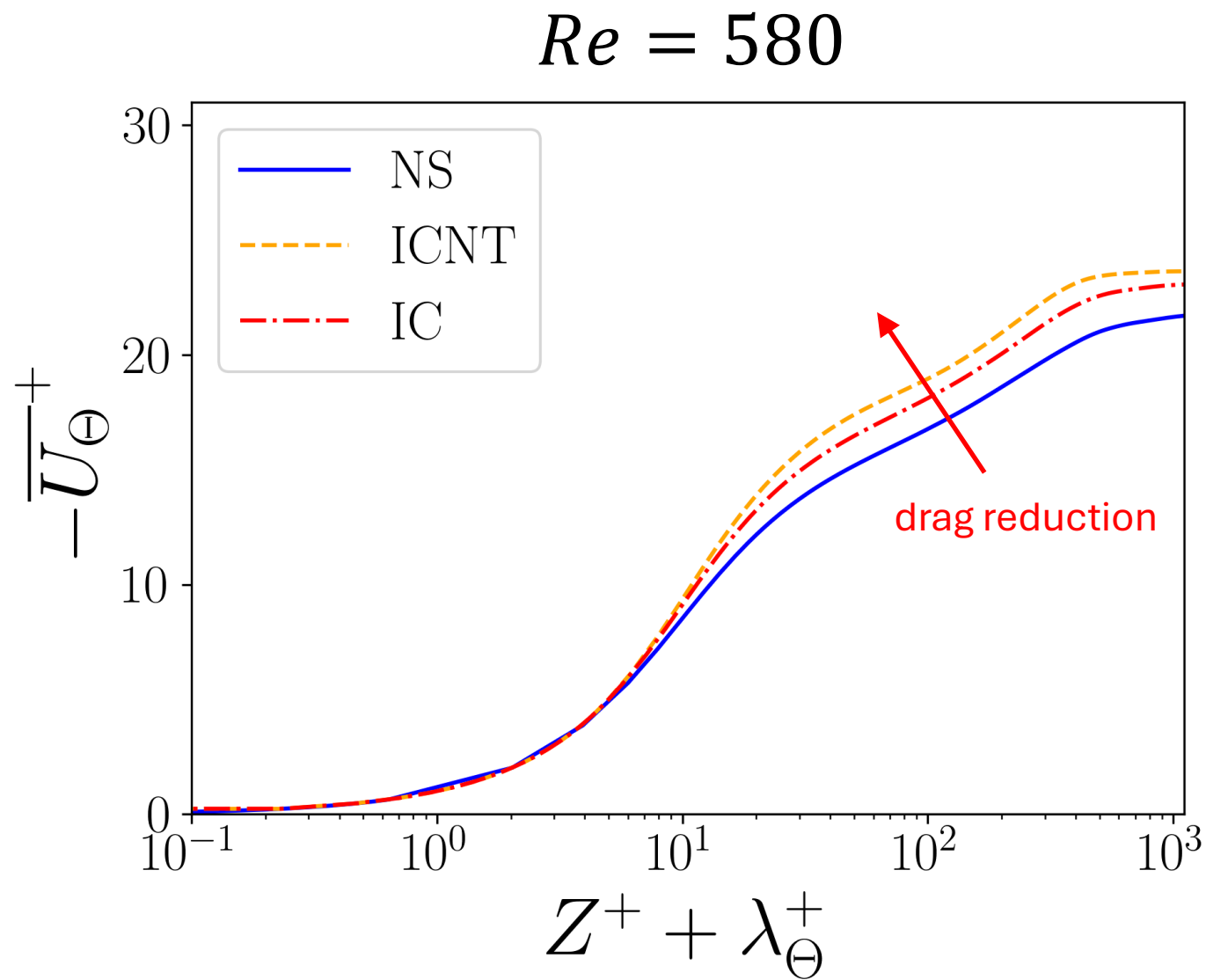
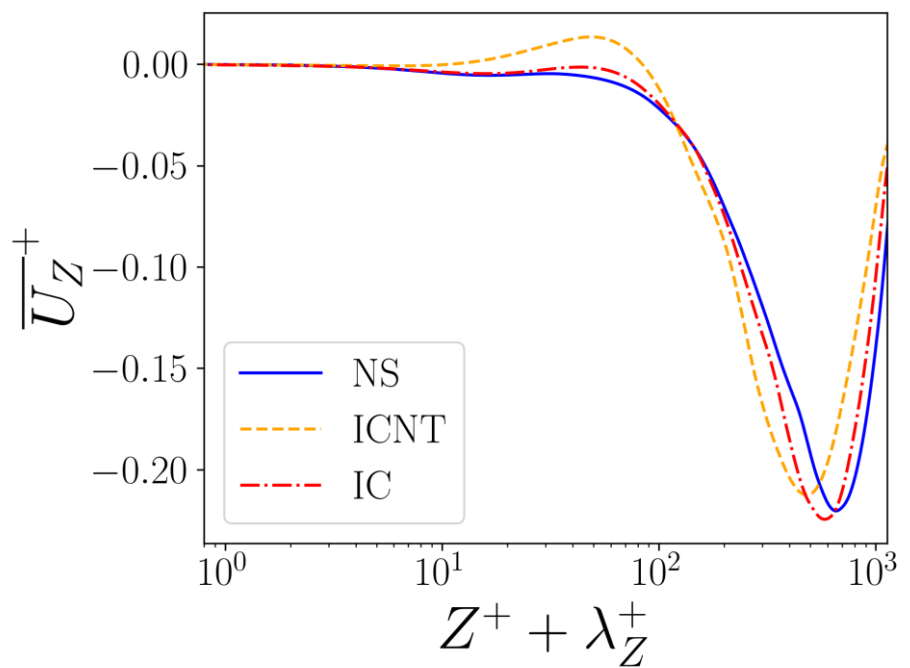
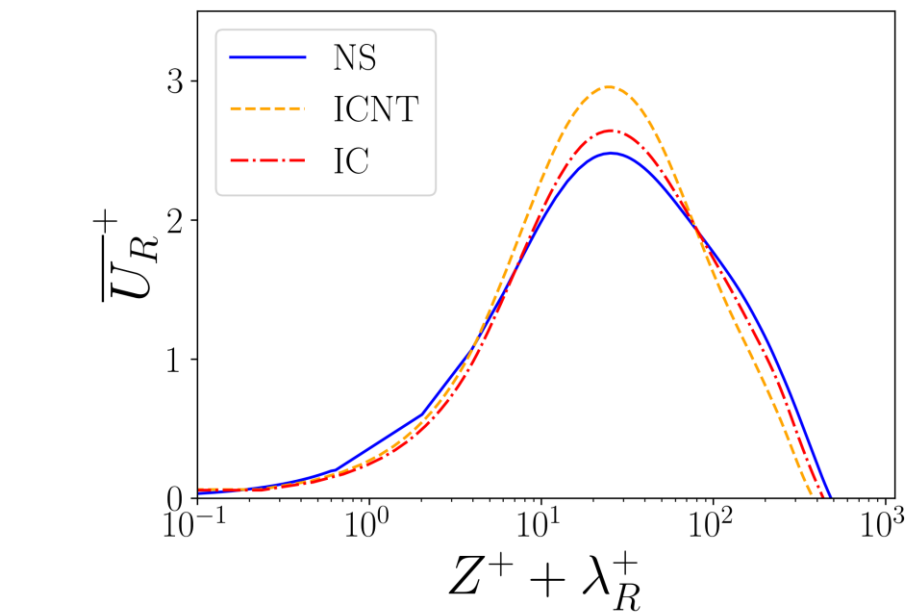




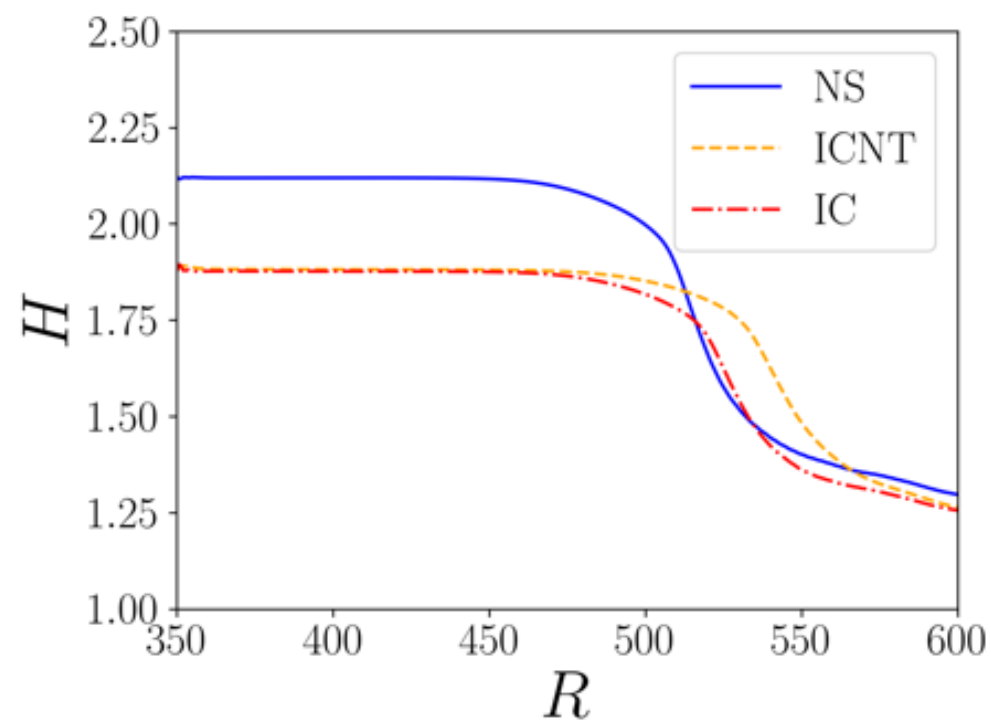
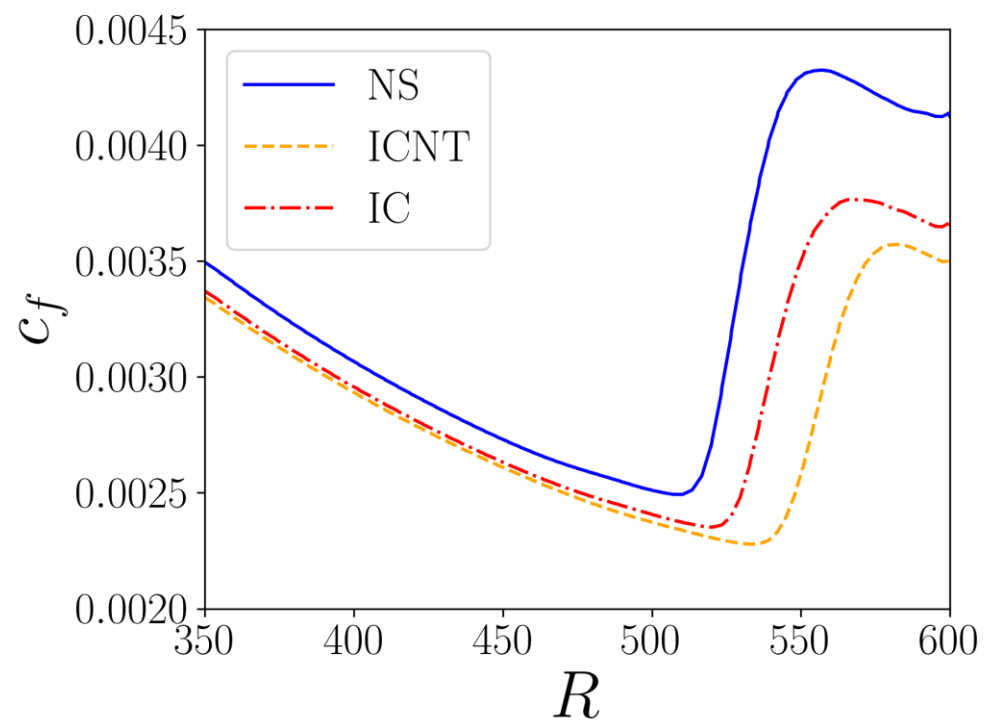


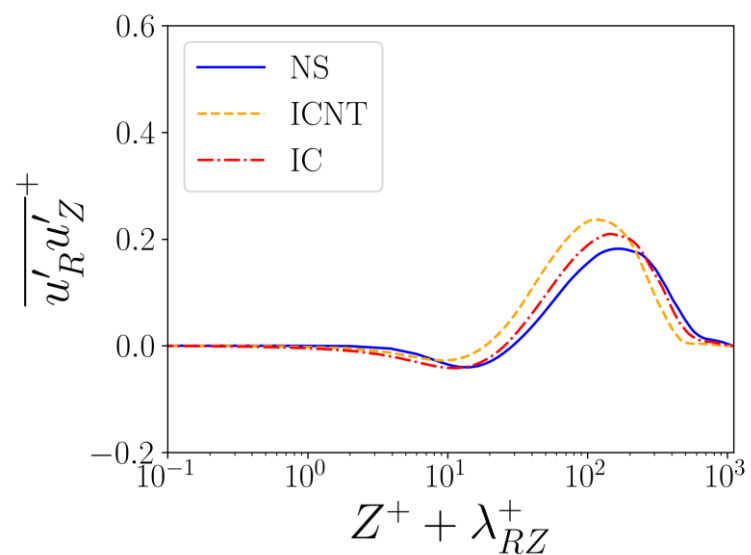
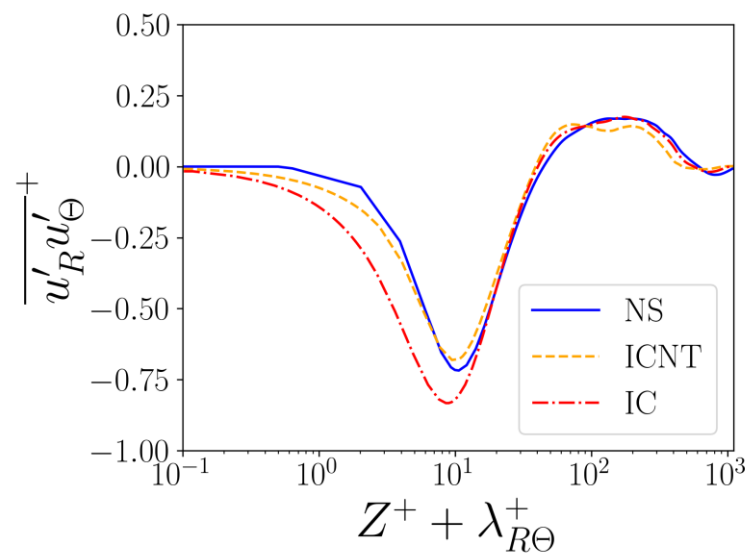




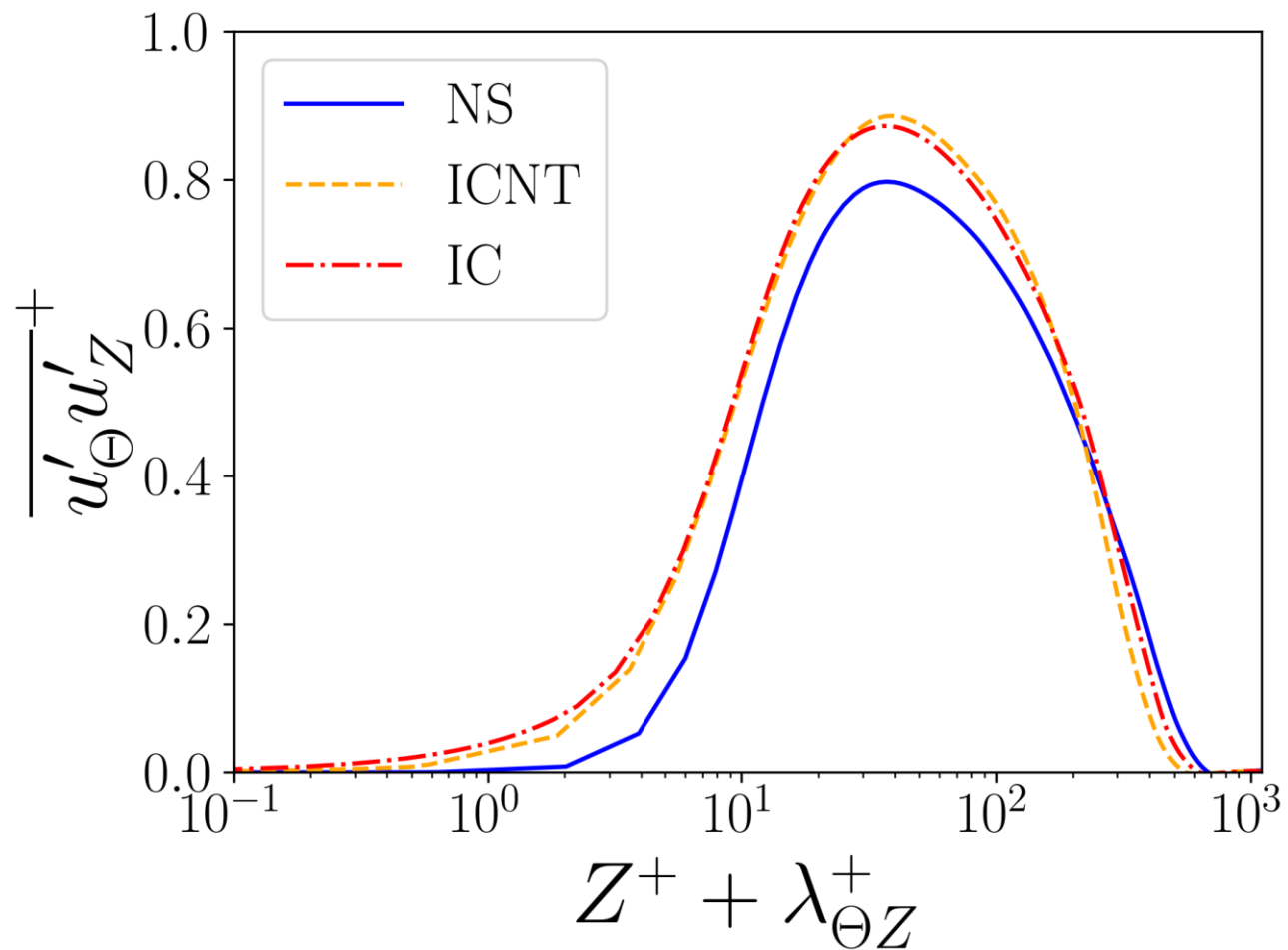




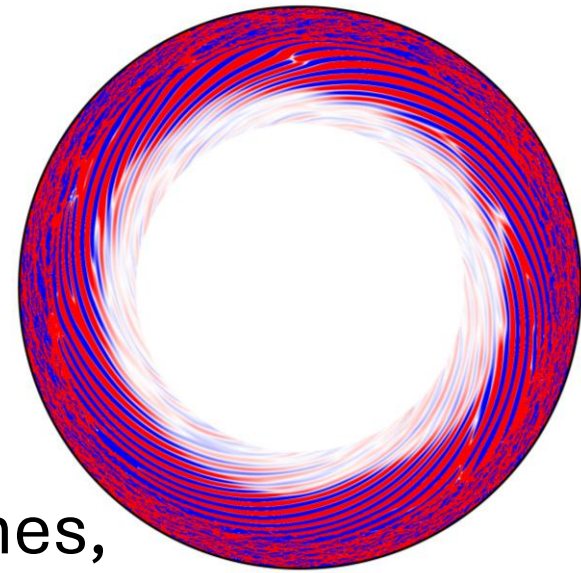




Case	$\lambda_R^+$	$\lambda_\Theta^+$	$\lambda_Z^+$	$\lambda_{R\Theta}^+$	$\lambda_{RZ}^+$	$\lambda_{\Theta Z}^+$
NS	0.000	0.000	0.000	0.000	0.000	0.000
ICNT	4.377	4.377	0.000	2.184	0.000	0.000
IC	4.483	4.483	3.012	2.242	1.780	1.805



# Conclusions



- **Homogenization** is an effective tool to model regularly microstructured rough wall, for quick parametric searches, particularly when coupled to **LSA** which provides immediate and reliable answers as to the performance of the surface
- Slip alone is insufficient to model the behavior of rough walls under transitional/turbulent conditions; a **transpiration** velocity at the fictitious  $Z = 0$  boundary is indispensable
- Isotropically arranged cones on the surface, periodicity 0.2 b.l. thicknesses, height 0.1 b.l. thicknesses, appear capable to slightly delay the breakdown to turbulence and reduce **skin friction**