



# BIOSKINS: Beyond Darcy's law in deformable porous media

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# Overview

- 1 Introduction
  - Process definition
  - Applications
- 2 Macroscopic approach
  - Volume averaging
  - Flow-structure coupling
- 3 Fluid-solid force
  - Inertia and topology
  - Unsteady flow in elastic porous media

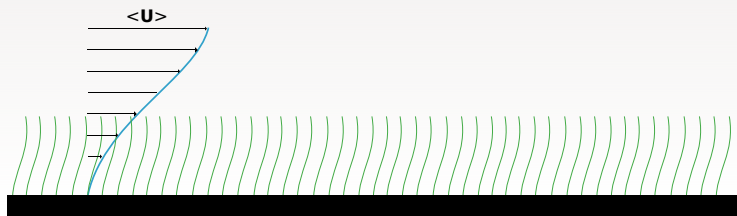


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# Typical poro-elastic layer



## Parameters:

- rigidity: restoring force vs hydrodynamic load
- sparsity
- macroscopic Reynolds number:  $\frac{hU}{\nu}$
- ...

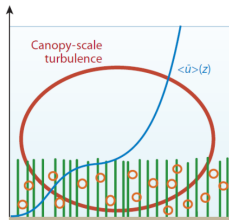
For a certain range of parameters we observe a regime of large scale turbulence (honami).  
Our final goal is to capture this large scale turbulence through homogenized model.





# Regime of large scale turbulence

- flow instability
- flow-structure coupling, coherent eddies
- $l_{micro}$ ,  $L_{macro}$



$$\begin{cases} u = \bar{u} + u' \\ v = \bar{v} + v' \end{cases} \quad (1)$$

Vertical momentum transfert

- $u' > 0, v' < 0$ : sweep
- $u' < 0, v' > 0$ : ejection

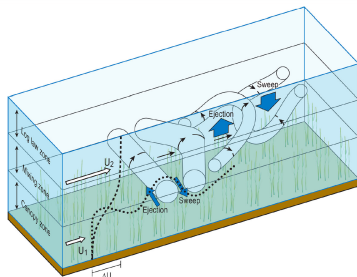


Figure: 3D aspect.

Ref: [Ghisalberti and Nepf, 2002], [Marjoribanks et al., 2017]



# Applications

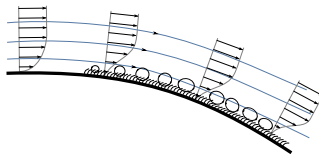
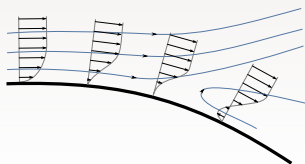


Figure: Innovative coating.



Figure: Field damage.

- sediment and nutrient transport
- vegetation management
- **passive flow control**

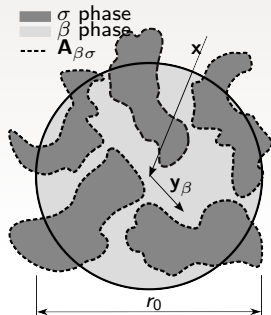


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# Macroscopic approach



Volume Averaged Navier-Stokes equations [VANS]

$$\frac{\partial \langle \mathbf{v}_\beta \rangle^\beta}{\partial t} + \langle \mathbf{v}_\beta \rangle^\beta \cdot \nabla \langle \mathbf{v}_\beta \rangle^\beta = - \frac{1}{\rho_\beta} \nabla \langle p_\beta \rangle^\beta + \nabla \cdot \frac{\mu_\beta}{\rho_\beta} \left[ \nabla \langle \mathbf{v}_\beta \rangle^\beta + \mathbf{T} \cdot \right]$$

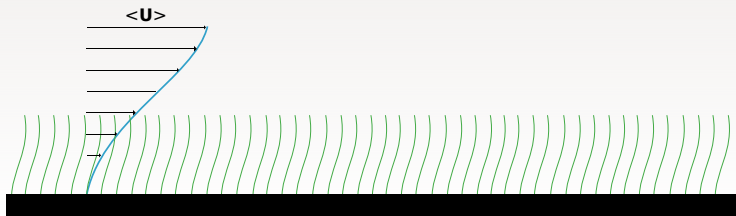
$$+ \frac{1}{\rho_\beta} \underbrace{\mathbf{D}_{\beta\sigma}}_{\text{fluid-solid force}} - \nabla \cdot \underbrace{\langle \tilde{\mathbf{v}}_\beta \tilde{\mathbf{v}}_\beta \rangle^\beta}_{\text{subgrid scale stresses}},$$

$$\nabla \cdot \langle \mathbf{v}_\beta \rangle^\beta = 0.$$

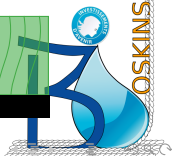
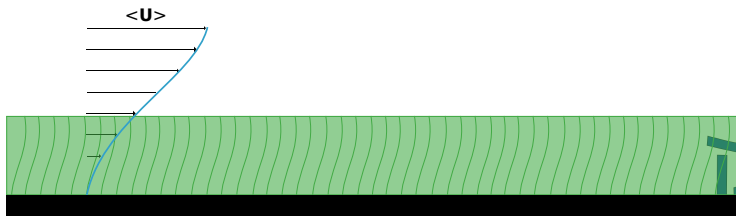
Refs: [Quintard and Whitaker, 1994],[Davit and Quintard, 2017]



# Homogenization of the sparse solid phase



↓ Rheology?



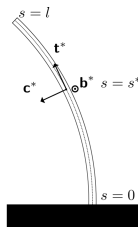
# Simplest beam mechanics

- $\rho_\sigma$  density of the solid
- $B$  bending rigidity
- $I_{xx}$  second moment of stem area

[EULER]

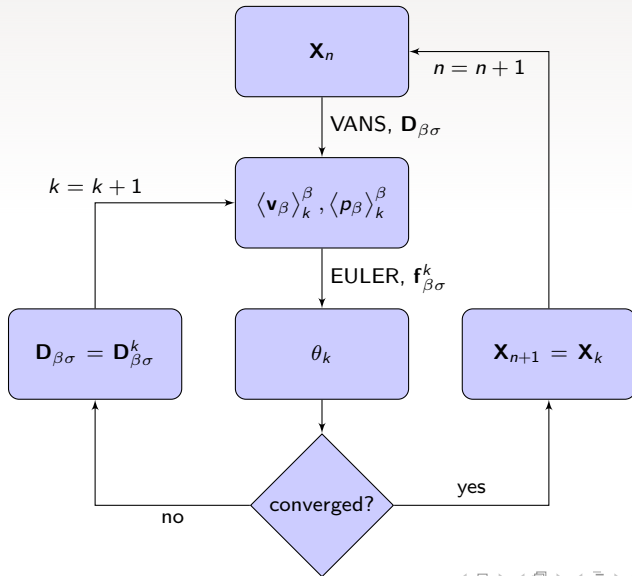
$$I_{xx} \rho_\sigma \frac{d^2\theta}{dt^2} \Big|_{s=s^*} = \frac{d}{ds} \left( B \frac{d\theta}{ds} \right) \Big|_{s=s^*} + \mathbf{c}^* \cdot \int_{s^*}^l \mathbf{f}_{\beta\sigma} ds \quad (3)$$

solver with variable time step for versatility



# Macroscopic code

Let  $\mathbf{X} = (\langle \mathbf{v}_\beta \rangle^\beta, \langle p_\beta \rangle^\beta, \theta)$ .  $\mathbf{D}_{\beta\sigma}$  and  $\mathbf{f}_{\beta\sigma}$  are function of  $\mathbf{X}$ .



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# From Darcy's law to inertial, steady regime

$$\begin{cases} \langle \mathbf{v}_\beta \rangle^\beta = v \boldsymbol{\lambda} \\ \mathbf{D}_{\beta\sigma} = D_\lambda \boldsymbol{\lambda} + \mathbf{D}_\perp \end{cases}, \text{ we write } D_\lambda = -\epsilon \mu K_D^{-1} (1 + F_\lambda) v \quad (4)$$

$K_D$  is a scalar permeability  $K_D^{-1} = \boldsymbol{\lambda} \cdot \mathbf{K}_D^{-1} \cdot \boldsymbol{\lambda}$ .

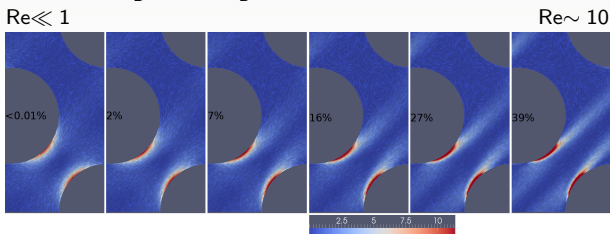


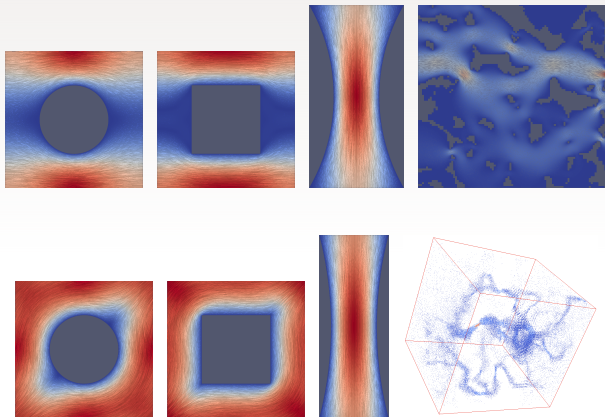
Figure: Normalized viscous dissipation

$\ell = ?$

$$Re = \frac{\ell v}{\nu}$$



# From Darcy's law to inertial, steady regime: geometries



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short-cut

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medium

DSA

disks in staggered arrangement

DRA

disks in regular arrangement

SSA

squares in staggered arrangement

SRA

squares in regular arrangement

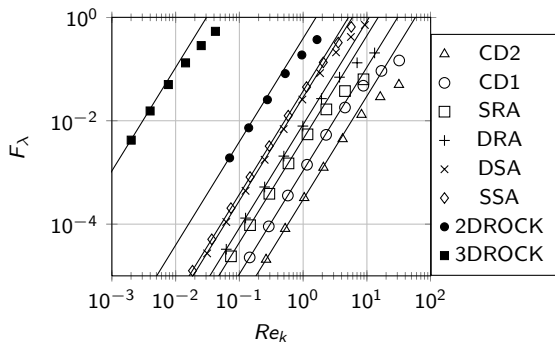
CD1, CD2

convergent-divergent



# Is $Re_k$ the right scaling?

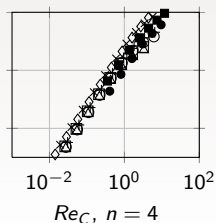
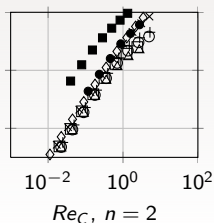
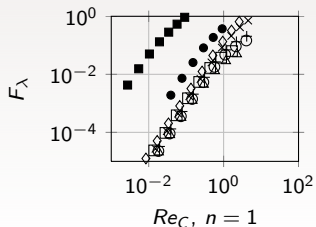
short-cut	medium
DSA	disks in staggered arrangement
DRA	disks in regular arrangement
SSA	squares in staggered arrangement
SRA	squares in regular arrangement
CD1, CD2	convergent-divergent



$$Re_k = \frac{v}{\nu} \sqrt{\frac{K_D}{\epsilon_\beta}} \quad (6)$$



# New scaling



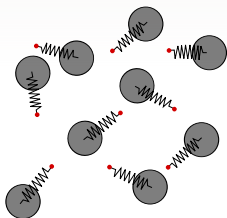
$$Re_C = \frac{\nu}{\nu} \sqrt{\frac{K_D}{\epsilon_\beta}} C_\lambda, \text{ with } \begin{cases} C_\lambda = \zeta(\mathbf{v}_0^* \cdot \nabla^* \mathbf{v}_0^*) \\ \zeta(\text{field}) = (\langle \|\text{field}\|^n \rangle^\beta)^{1/n} \end{cases} \quad (7)$$

medium	$C_\lambda$
CD2	0.10
CD1	0.19
SRA	0.36
DRA	0.49
DSA	0.62
SSA	0.77
2DROCK	5.77
3DROCK	176

[Pauthenet et al., under publication]



# Pore scale fluid-structure interaction



Refs: [Bigot et al., 2014]

## Parameters

- $m^* = \frac{\rho_\sigma}{\rho_\beta}$
- $f^* = \frac{\sqrt{k/m_\sigma}}{U/d}$
- $Re_d = \frac{dU}{\nu}$

We add two parameters in the dataset for the fluid-solid force  $\mathbf{D}_{\beta\sigma}$ .

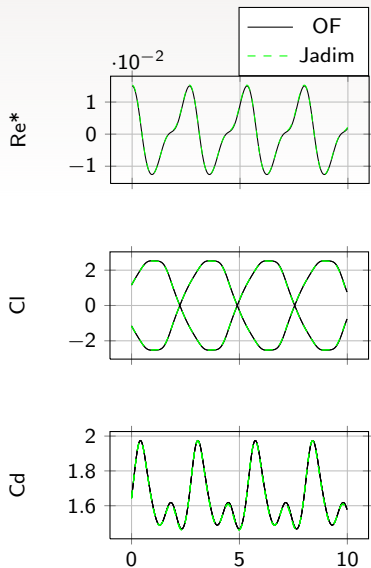
## Numerical method

- large displacement → immersed boundary method
- parallellized code

→ JADIM



# Rigid case



Comparison with OpenFOAM (OF).

$$Re = \frac{dV}{\nu} \sim 100$$

$$\begin{cases} Re^* = \frac{Re - \langle Re \rangle}{\langle Re_{ref} \rangle_f} \\ (C_d, C_l) = \frac{1}{\frac{1}{2} \rho \beta S U^2} \end{cases} \quad (8)$$



# Validation and goals

## Validation cases

- Periodic cell with elastic cylinders (different IBM)
- Early stage of wake behind a cylinder (experiments) [Bouard and Coutanceau, 1980]

## Goals

- $\mathbf{D}_{\beta\sigma}(Re_d, m^*, f^*)$  from DNS
- $l_{macro} \gg l_{micro}$ ?
- $\tau_{macro} \gg \tau_{micro}$ ?
- size of a representative volume

# Conclusion

## Key points

### Macroscopic code

- coupling of the VANS with EULER
- fluid-solid force from a dataset
- subgrid-scale stresses?

### Structure

- no rheology law
- simplest individual beam mechanics
- reduce the number of degrees of freedom?

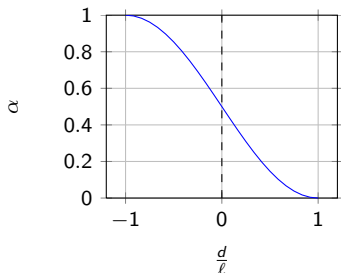
### Fluid-solid force

- weak inertia regime
- unsteady flow in elastic media
- implement simple solid-solid contacts?



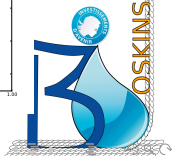
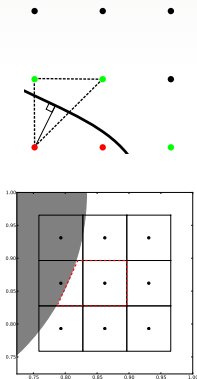
# Immersed Boundary Method: direct forcing


- $\mathbf{f}_\sigma$  imposes velocity in solid phase
- Pseudo Poisson equation for mass conservation
- Stress on solid obstacle from  $\mathbf{f}_\sigma$



a different method (Ghost cell + cut-cell)

- Dirichlet for predicted velocity
- Neuman for pressure correction



 Bigot, B., Bonometti, T., Lacaze, L., and Thual, O. (2014).

A simple immersed-boundary method for solid-fluid interaction in constant-and stratified-density flows.

*Computers & Fluids*, 97:126–142.



Bouard, R. and Coutanceau, M. (1980).

The early stage of development of the wake behind an impulsively started cylinder for  $40 < re < 10^4$ .

*Journal of Fluid Mechanics*, 101(3):583–607.



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Technical notes on volume averaging in porous media I: How to choose a spatial averaging operator for periodic and quasiperiodic structures.

*Transport in Porous Media*, pages 1–30.



Ghisalberti, M. and Nepf, H. (2002).

Mixing layers and coherent structures in vegetated aquatic flows.

*Journal of Geophysical Research*, 107.



Marjoribanks, T. I., Hardy, R. J., Lane, S. N., and Parsons, D. R. (2017).

Does the canopy mixing layer model apply to highly flexible aquatic vegetation? Insights from numerical modelling.

*Environmental Fluid Mechanics*, 17(2):277–301.



Quintard, M. and Whitaker, S. (1994).

Transport in ordered and disordered porous media II: Generalized volume averaging.

*Chemical Engineering Science*, (14,15).

